

TUNING FORKS AS A GENERATOR OF QUANTUM TURBULENCE IN SUPERFLUID HELIUM

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In honor of Evgenii Kuznetsov’s 70th birthday

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Report outline

- Quantum turbulence in superfluid helium
- Vortex generation by vibrating bodies
- Tuning forks, electro-mechanical coefficient
- Tuning fork as a turbulence generator
- Tuning fork as a vortex detector
- Two regimes of vortex generation
- New type of tuning forks

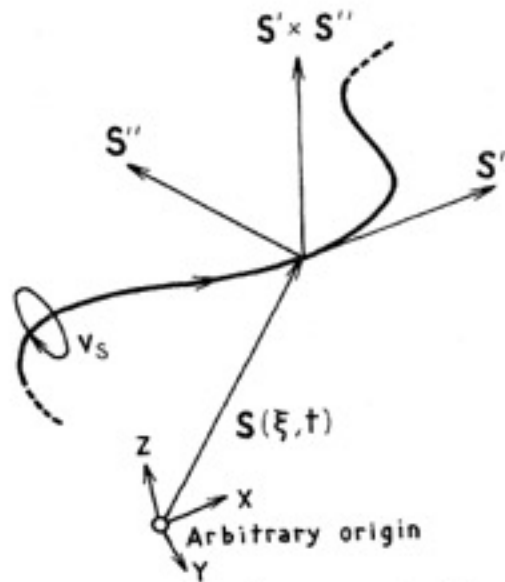


FIG. 1. Space curve representing a vortex line with its position described as $\mathbf{S}(\xi, t)$, when ξ is arclength. $\mathbf{S}' = d\mathbf{S}/d\xi$ is a unit vector along the vortex line; $\mathbf{S}'' = d^2\mathbf{S}/d\xi^2$ is the local curvature vector (whose magnitude is $1/R$); and $\mathbf{S}' \times \mathbf{S}''$ is binormal, which also has the magnitude $1/R$ (Donnelly, 1991, Fig. 1.14).

A motion of superfluid component could be written as

$$\boldsymbol{\omega}(\mathbf{r}) = \text{rot} \mathbf{v}_s = \kappa \int d\mathbf{S} \delta(\mathbf{r} - \mathbf{S}(\xi, t))$$

where the integration is along the line singularity $\mathbf{S}(\xi, t)$.

A single infinite straight line induces a velocity field, described in cylindrical coordinates by

$$\mathbf{v}_s = \left[0, \frac{\kappa}{2\pi r}, 0 \right]$$

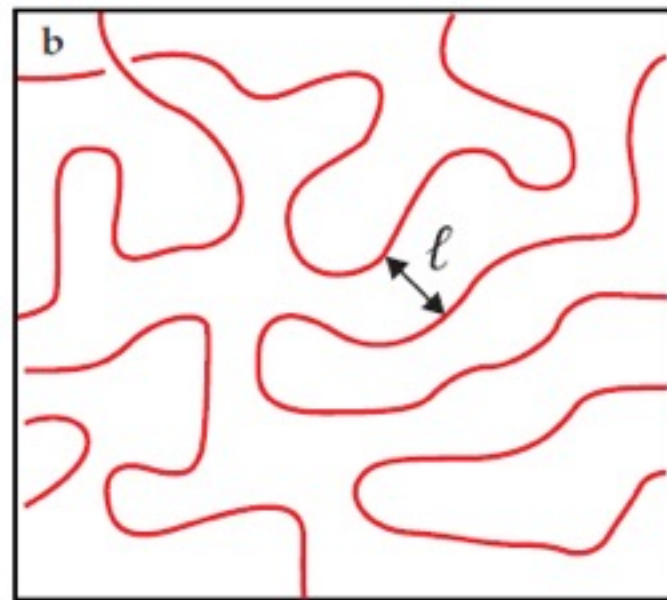
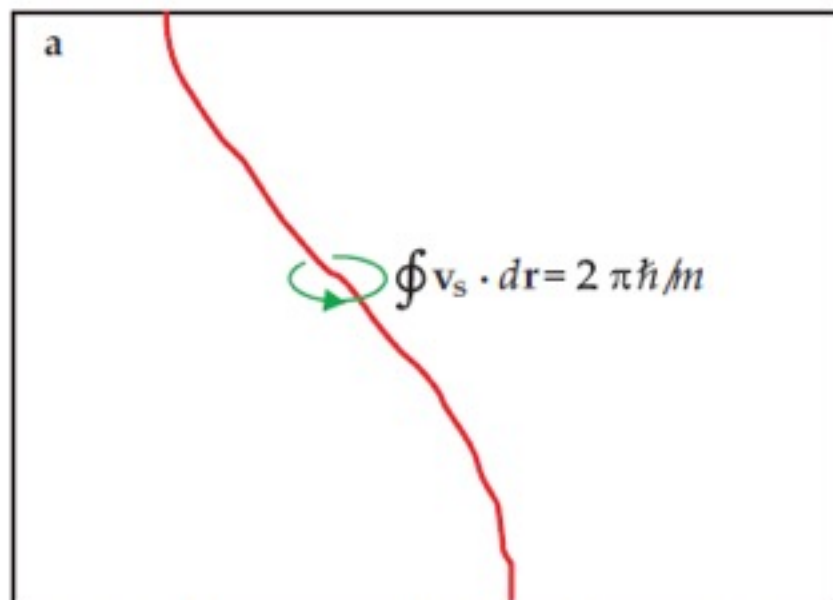
Having only an azimuthal component which increases rapidly on approaching the vortex axis. In particular, at $r \leq 3\text{\AA}$, the velocity v_s becomes greater than the velocity of roton generation, i.e., Landau's velocity.

Vortexes in superfluid helium

Quantization of vortex lines, core radius a_0 and intervortex distance ℓ

$\oint \mathbf{v}_s \cdot d\mathbf{r} = n \kappa$, where $\kappa = \frac{2\pi\hbar}{M}$ is the circulation quant.

$M = 4$ for ^4He and $M = 6$ for a pair of ^3He atoms.



- ℓ is the mean intervortex distance,
- Vortex core radius $a_0 \simeq 1 \text{ \AA}$ for ^4He & $a_0 \simeq 800 \text{ \AA}$ at low p .
for ^3He

Classical turbulence – measures of its intensity

Reynolds number

For isothermal flows

T(P)

$$Re = \frac{UL}{\nu}$$

ν (cm²/s)=10⁻² St

kinematic viscosity

air	20 C	0,15
water	20 C	1,004x10 ⁻²
ethanol	20 C	0.022
mercury	20 C	1,2x10 ⁻³
Helium I	2,25 K (VP)	1,96x10 ⁻⁴
Helium II	1,8 K (VP)	9,01x10 ⁻⁵
He-gas	5,5 K (2,8 bar)	3,21x10 ⁻⁴

•He II and ³He B – so far the only two media where quantum turbulence has been systematically studied under controlled laboratory conditions

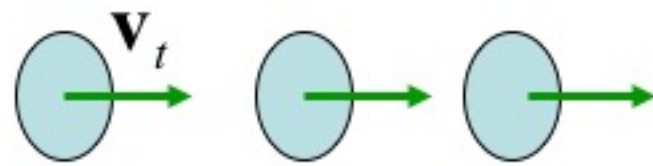
Counterflow turbulence phenomenology (Vinen 1957)

Vortex ring

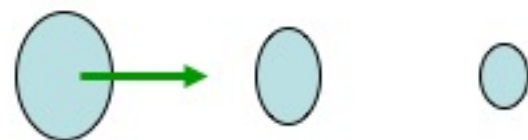


$$v_t = \frac{\kappa}{4\pi b} \left(\ln \frac{8b}{a} - \frac{1}{4} \right) \cong \frac{\kappa}{b}$$

$T \rightarrow 0$



Finite T



In counterflow, though, if $|v_t| < |v_n - v_s| = V_{CF}$ rings with $b > b_c$ expand

Dimensional analysis and analogy with classical fluid dynamics leads to the Vinen equation:

$$\frac{dL}{dt} = \underbrace{\chi_1 \frac{B}{2} \frac{\rho_n}{\rho} V_{CF} L^{3/2}}_{\text{production}} - \underbrace{\chi_2 \frac{\hbar}{m_4} L^2}_{\text{decay}} + (g(V_{CF}))$$



Reproduced by Schwarz (1988) –
computer simulations
Local induction approximation
Importance of reconnections

^4He



T

normal liquid He I

Classical Navier-Stokes fluid of extremely low kinematic viscosity

2



Superfluid transition at $T_c=2.17$ K

He II –

a “mixture” of two fluids

normal fluid of extremely low kinematic viscosity

+

Inviscid superfluid

3

Circulation is quantized

$$\kappa = \frac{2\pi\hbar}{m_4} \approx 10^{-3} [\text{cm}^2 / \text{s}]$$



T \rightarrow 0 limit

1

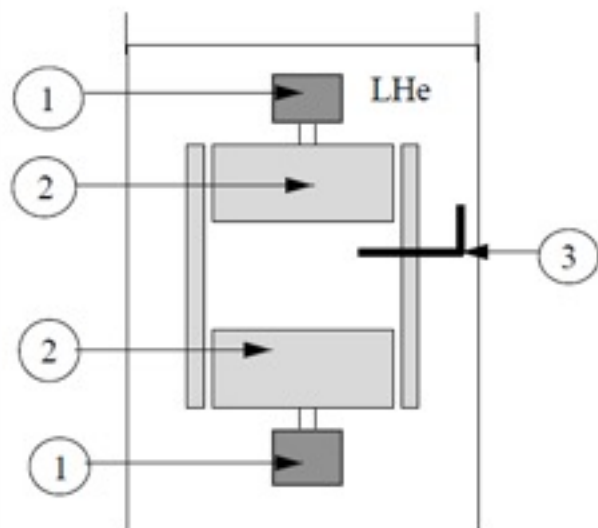
0

Pure superfluid

Maurer, Tabeling,

Europhysics Lett. 43 (1998) 29

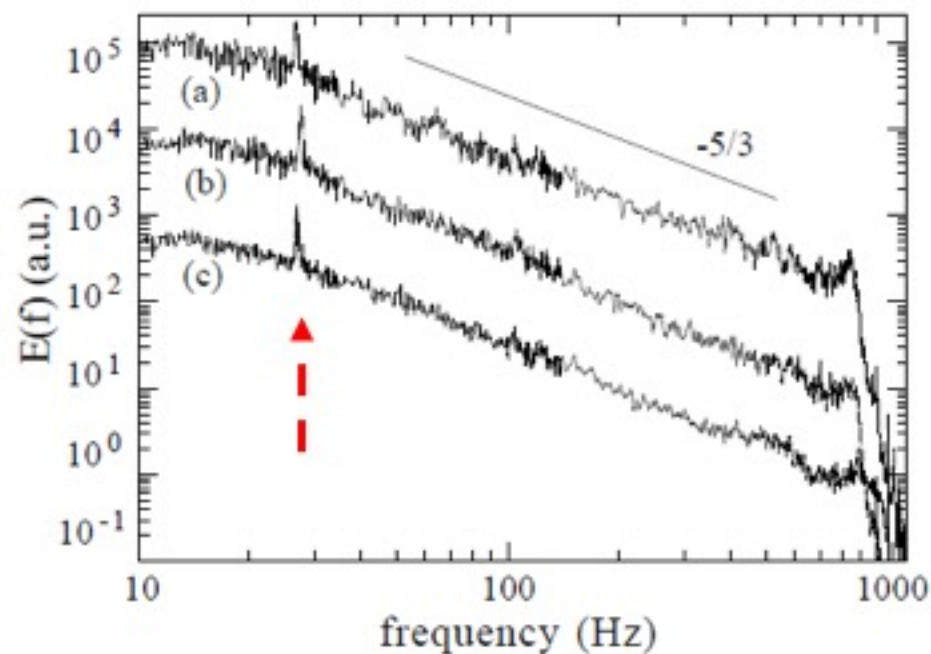
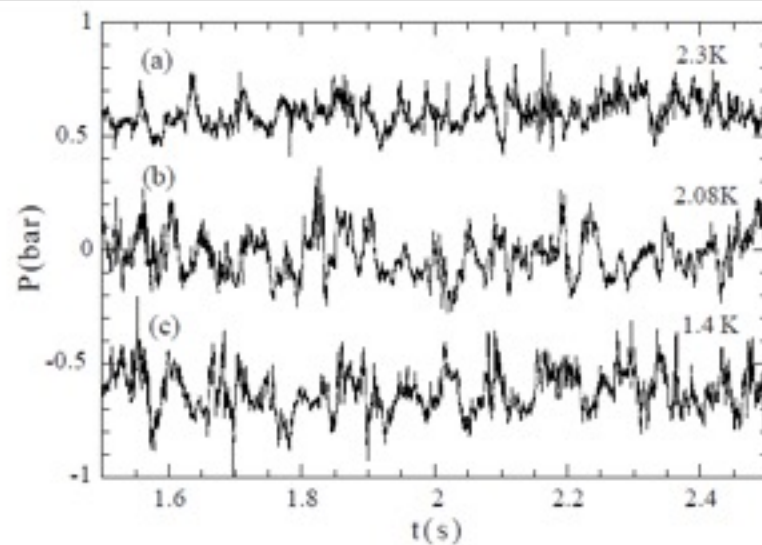
Flow between counterrotating discs



Energy spectra obtained in the same conditions:

$U=80\text{cm/s}$; $Re=2\times 10^6$ a-2.3K; b-2.08K; c-1.4 K

The spectra have been shifted vertically so as to make their representation clear.

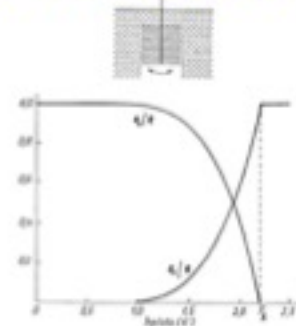


Oscillating objects used in experiments in He II and in ^3He

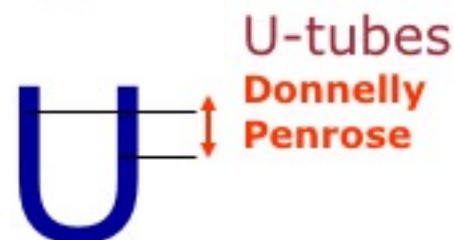
Discs and piles of discs



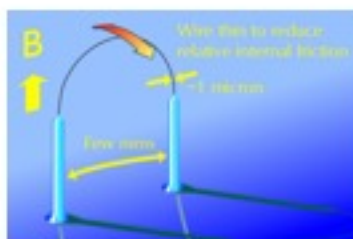
Torsional oscillators



Andronikashvili



Wires
He II and
 ^3He



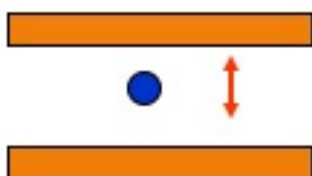
Many authors – vibrating wire viscometers

Vinen

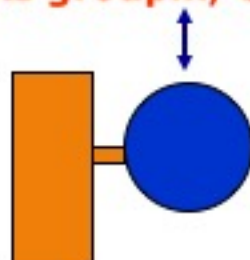
Morishita, Kuroda, Sawada, Satoh, JLTP 76, 387 (1989)

Lancaster – Pickett's group, Osaka – Yano et al., Kosice Skyba et al., Moscow Dmitriev et al., Helsinki- YKI group..., Grenoble Bunkov et al., ...

Spheres
He II
(^3He ?)



Schoepe et al



Luzuriaga

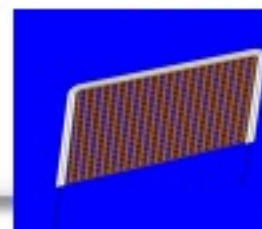
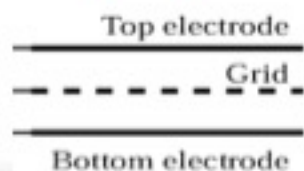


Hollis-Hallett

Lancaster – P. McClintock's group

Lancaster – G. Pickett's group

Grids
He II and
 ^3He



Typical characteristics of commercially available forks

Room temperature
closed fork



frequency 2^{15} Hz = 32764 Hz
linewidth around 0.5 Hz

Bare fork



frequency down by 7-10 Hz
linewidth around 5-6 Hz
linewidth around 3 Hz

Ground off can

Fork in vacuum: LN2 32722 Hz linewidth 0.2 Hz
LHe 32710 Hz linewidth 0.06 Hz
Q about 500 000

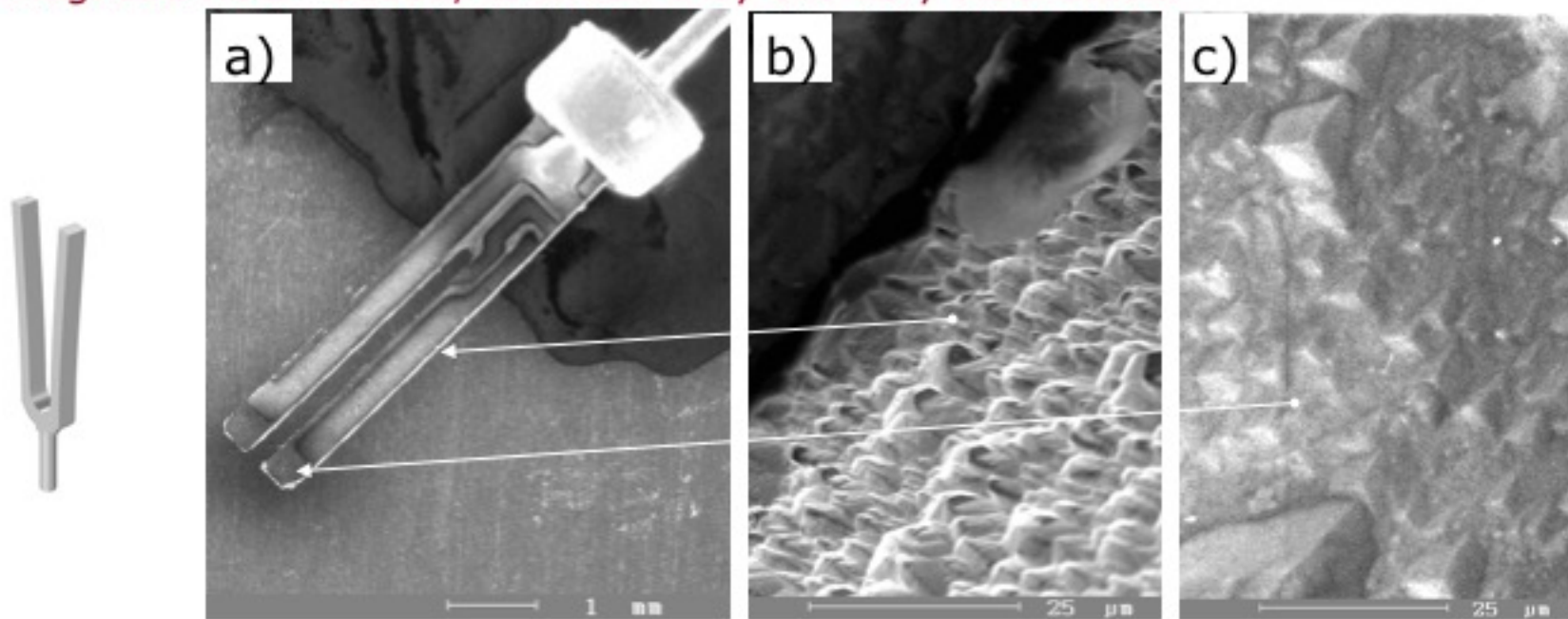
Fork in cryofluids: He I SVP 4.2 K 31600 - 31700 H
linewidth 12-13 Hz
He II SVP 1.3 K linewidth 1.2 Hz

Response to about 8 orders of magnitude of the ac drive (up to 130 V rms) fork rms velocity reaches up to 10 m/s

Thermal cycling is not a serious issue for most applications

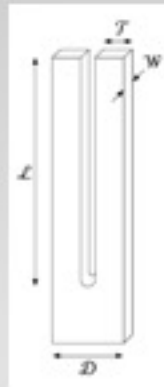
Quartz tuning forks –

- Commercially produced piezoelectric oscillators, used as frequency standards in watches (2^{15} Hz = 32 768 Hz at room temperature)
- New addition to a family of oscillating objects, a probe to investigate physical properties of cryogenic fluids, especially gaseous He, He I, He II and $^3\text{He-B}$
- Cheap, robust, widely available, easy to install and use (2 wires only), very sensitive
- No magnetic field needed, to which they are very insensitive



An electron micrograph of the quartz tuning fork (a) and details of its side (b) and top (c) quartz surface.

Quartz Tuning Fork: Thermometer, Pressure- and Viscometer for Helium Liquids, R. Blaauwgeers, M. Blazkova, M. Clovecko et al, *JLTP*, 146, 5/6, 537 (2007)



- **Mechanical properties**

$$\frac{d^2x}{dt^2} + \gamma \frac{dx}{dt} + \frac{k}{m}x = \frac{F}{m}$$

- Resonance frequency

$$\omega_0 = \sqrt{\frac{k}{m}}$$

- Quality $Q = \frac{\omega_0}{\gamma}$, where width of resonance curve $\Delta\omega = \gamma$

- Effective mass of prong

$$m_{\text{vac}} = 0.24267 \rho_q \mathcal{L} \mathcal{W} T$$

$$\begin{aligned} F_0 &= (a/2) U_0, \\ R &= 2m\gamma/a^2, \\ L &= 2m/a^2, \\ C &= a^2/(2k). \end{aligned}$$

- **Electrical properties**

- Fork response $I(t) = a \frac{dx(t)}{dt}$
 $U = U_0 \cos(\omega t)$

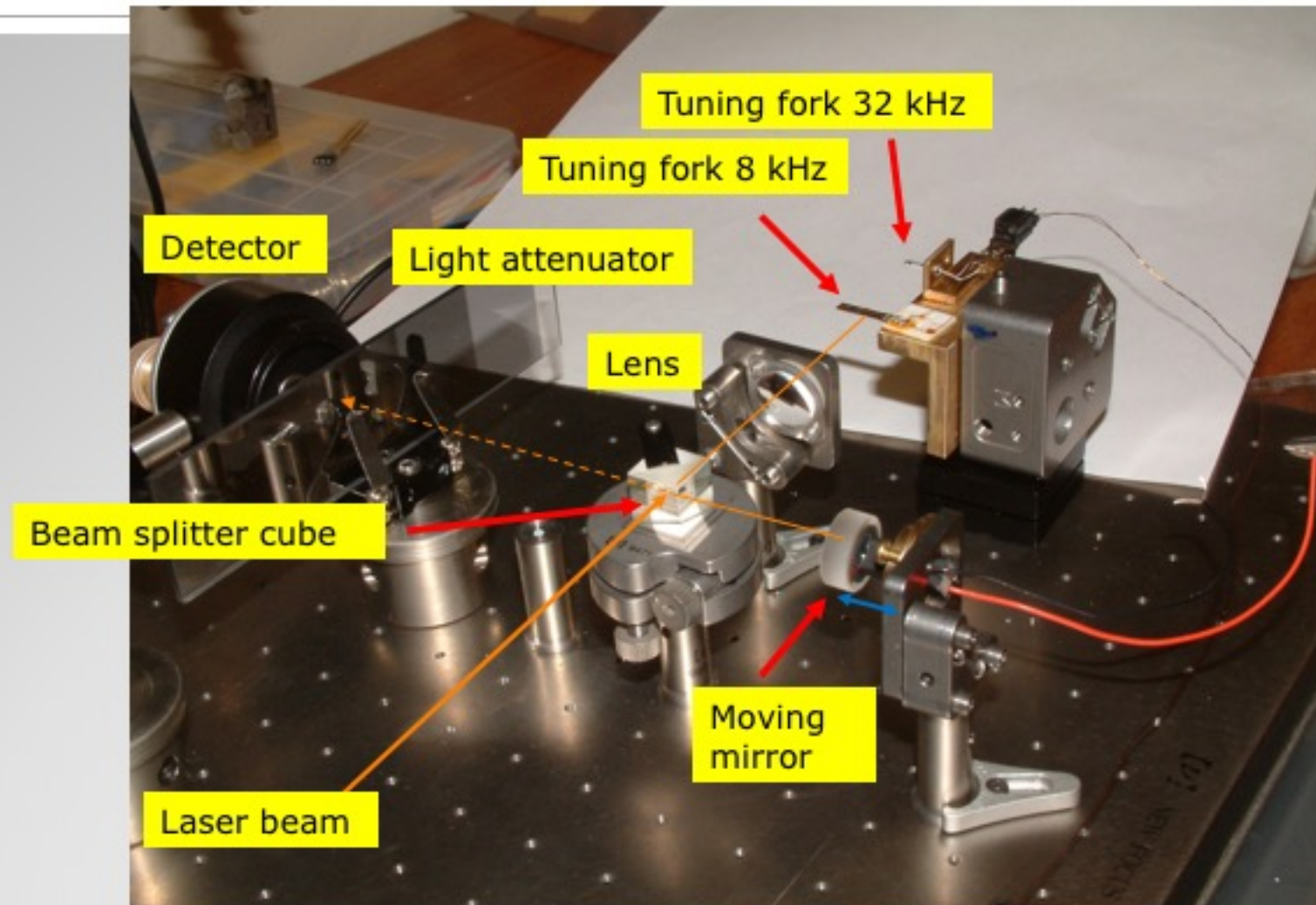
- The mechanical oscillator with the equivalent electrical RLC series resonance circuit. The corresponding differential equation for the current is

$$\frac{d^2I}{dt^2} + \frac{R}{L} \frac{dI}{dt} + \frac{1}{LC} I = \frac{1}{L} \frac{dU}{dt} \Rightarrow$$

$$\omega_0^2 = 1/(LC), \quad \gamma = R/L, \quad \text{and} \quad 1/L = (F_0/U_0) a/m$$

$$a = \sqrt{\frac{2m \Delta\omega}{R}}$$

Electro-mechanical calibration



Michelson interferometer scheme for measurement of amplitude of fork vibration

Absolute prong's velocity (T_{room} , air)

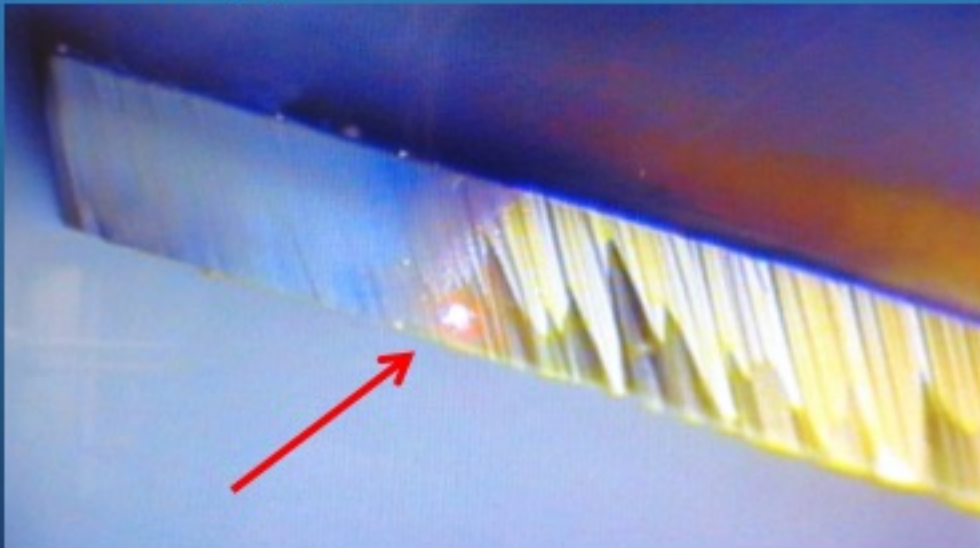
Bending vibration $f_c=75,64$ kHz

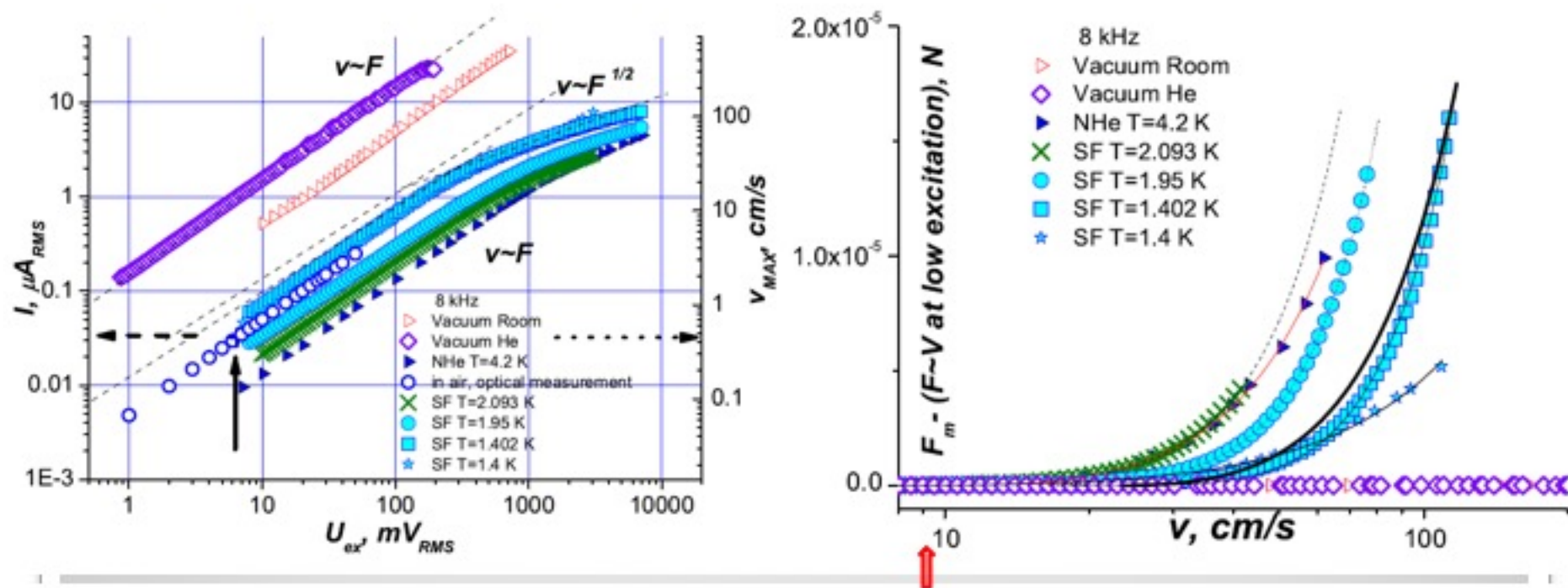
$V_{\text{drive}}=800$ mV_{pp}; $f = 76,228$ kHz, $\Delta f=17$ Hz;

sensitivity of Doppler analyzer (DA) 100 mm/s/V

signal of DA = 4.8 V, velocity of vibration $v_p=48$ cm/s

No changing in DA signal at scanning perpendicular to prong width. Laser point was near the end of the fork's prong ($v_{\text{max}} \sim 1.2 * v_p$)





At low excitation the fork response and corresponding velocity are proportional to the driving force. Higher excitation leads to the appearance of some additional mechanism of energy dissipation and to a deviation of the fork response (and its velocity) from linear dependence. Note that such deviations were not seen in vacuum (left graphic); they are usually attributed to the transition to turbulence above a critical velocity.

We estimated the deviation of the applied force (and counteracting fluid forces) by subtracting off the extrapolated the linear dependence at low excitation.

$$F_{\text{add}} = F(v) - C \times v$$

The results are shown in right graphic. The fluid friction force increases very rapidly after the critical velocity v_c is attained. We fitted the deviation by power law $F_{\text{add}} = C_2 \times (v - v_c)^N$ at different temperatures.

Low temperature behaviour of 8kHz fork

Generation of the quantum turbulence by tuning fork

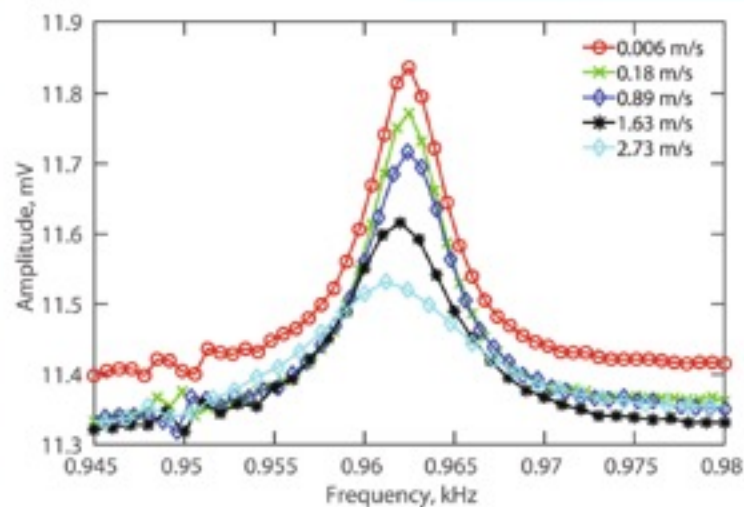
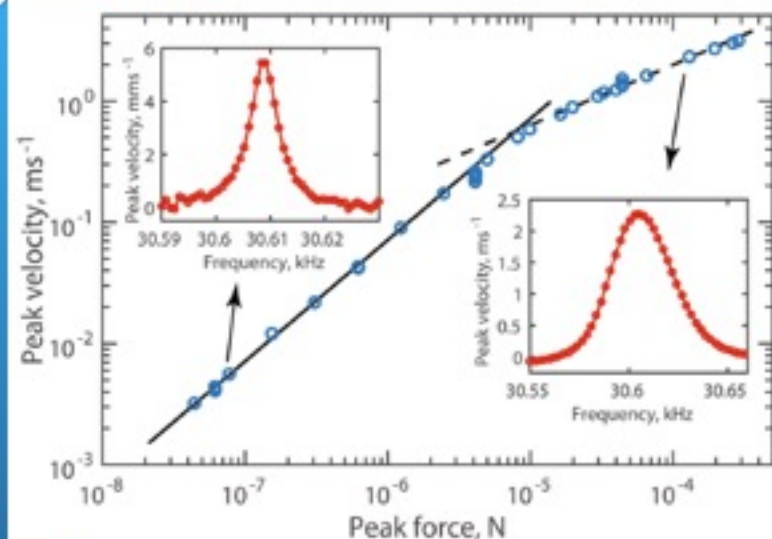
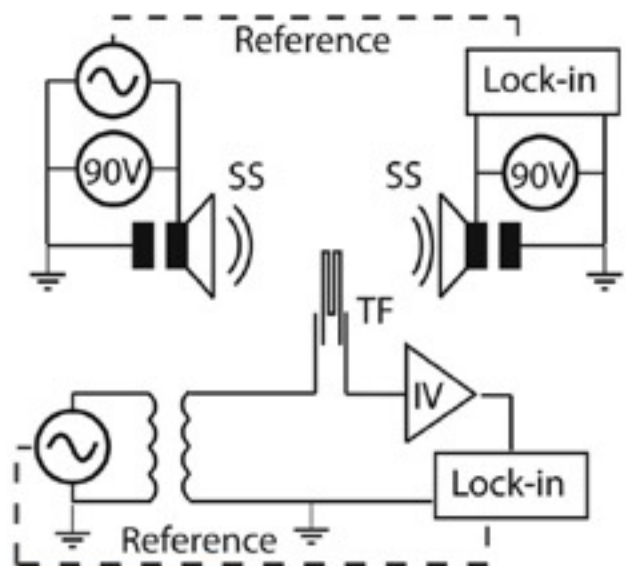


Fig. 3 Second sound frequency sweeps of the 5th second sound mode at various tuning fork velocities. The changes in the resonance frequency observed by the peaks in the second sound velocity due to the temperature drift of a few millikelvin and to the time increased damping (Color figure online)

M. J. Jackson et al., JLTIP, 183, 208 (2016)

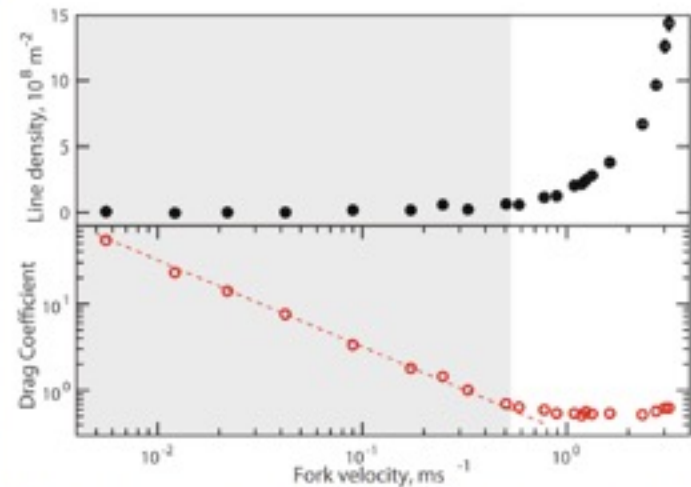
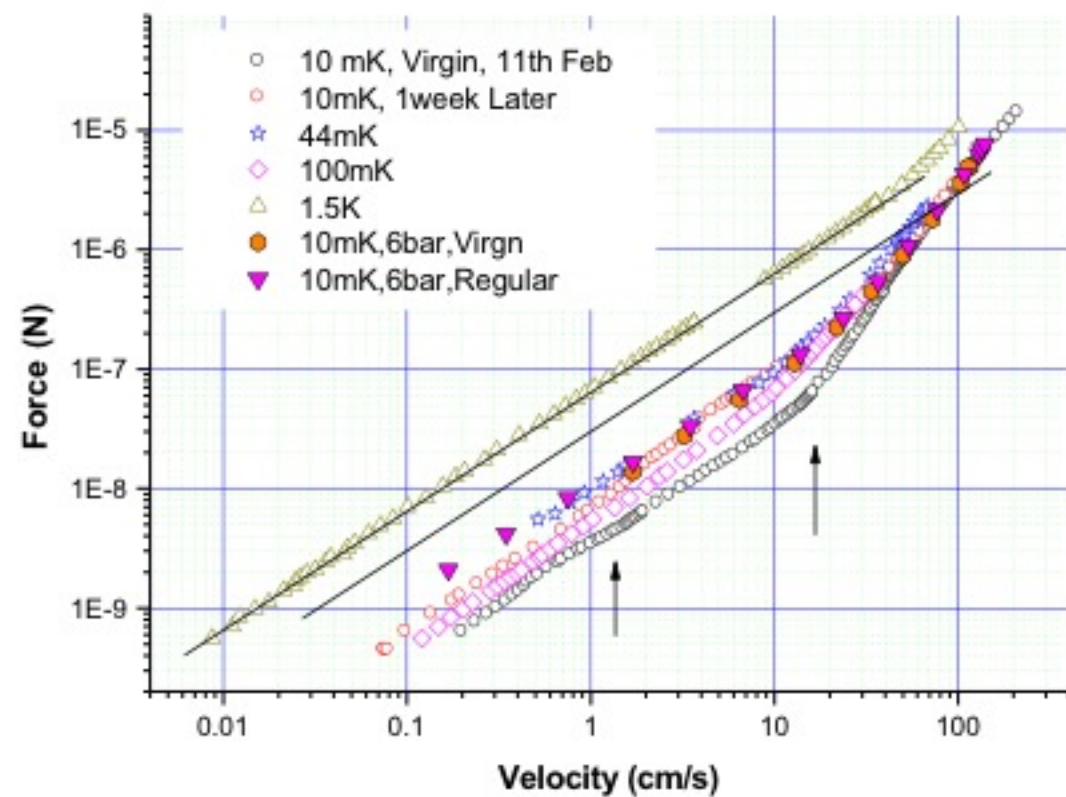


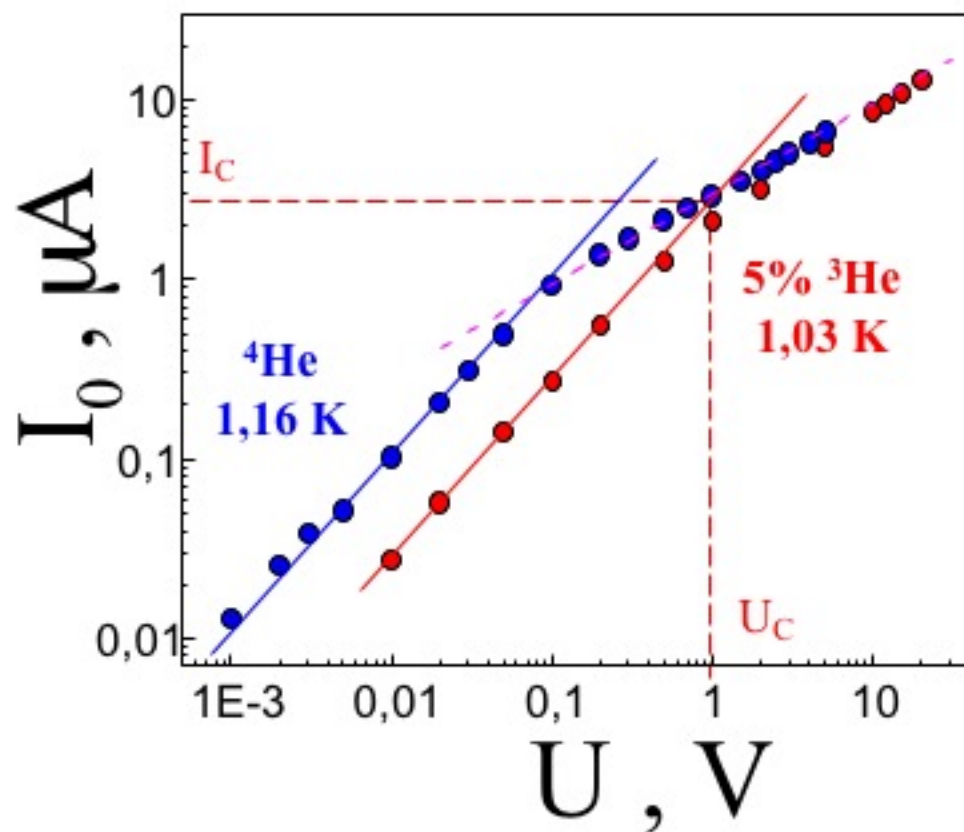
Fig. 4 Top Vortex line density inferred from the second sound attenuation as a function of tuning fork velocity. Bottom drag coefficient experienced by the tuning fork versus tuning fork velocity. The dashed line is a guide for the eye and the shaded area highlights the laminar regime. For details, see text (Color figure online)

Out Fork full results....all Temp , P = 5bar, 6bar



Temperature behaviour of fork response

Current-voltage characteristic ($T > 1$ K).



At low U : $U \sim I_0$

laminar regime

(solid lines)

At high U : $U \sim I_0^2$

turbulence

(dashed line)

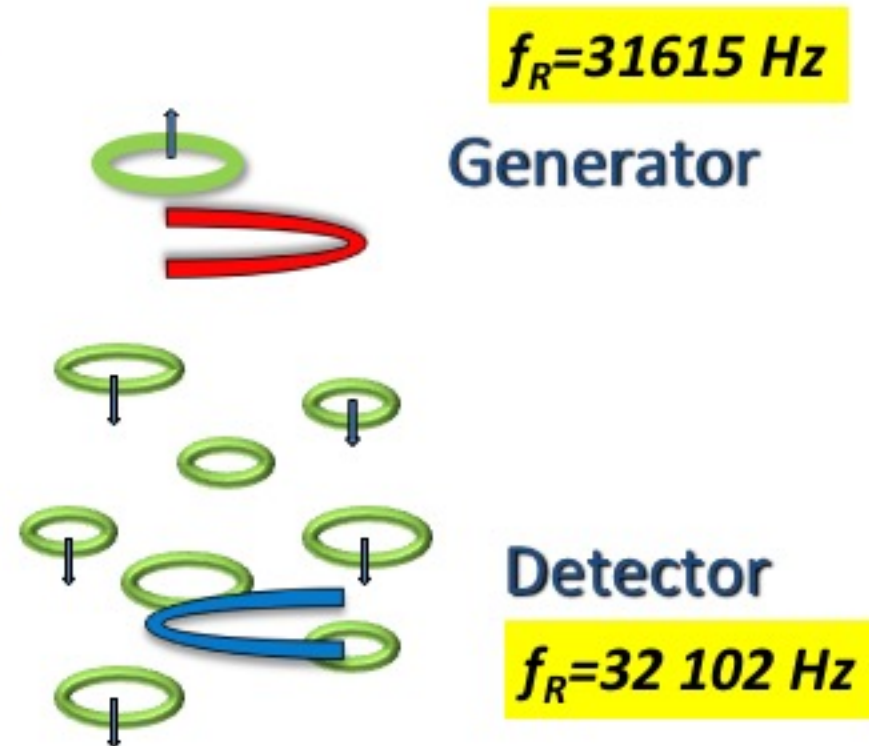
The laminar flow is more stable in the ^3He - ^4He mixture than that in pure ^4He

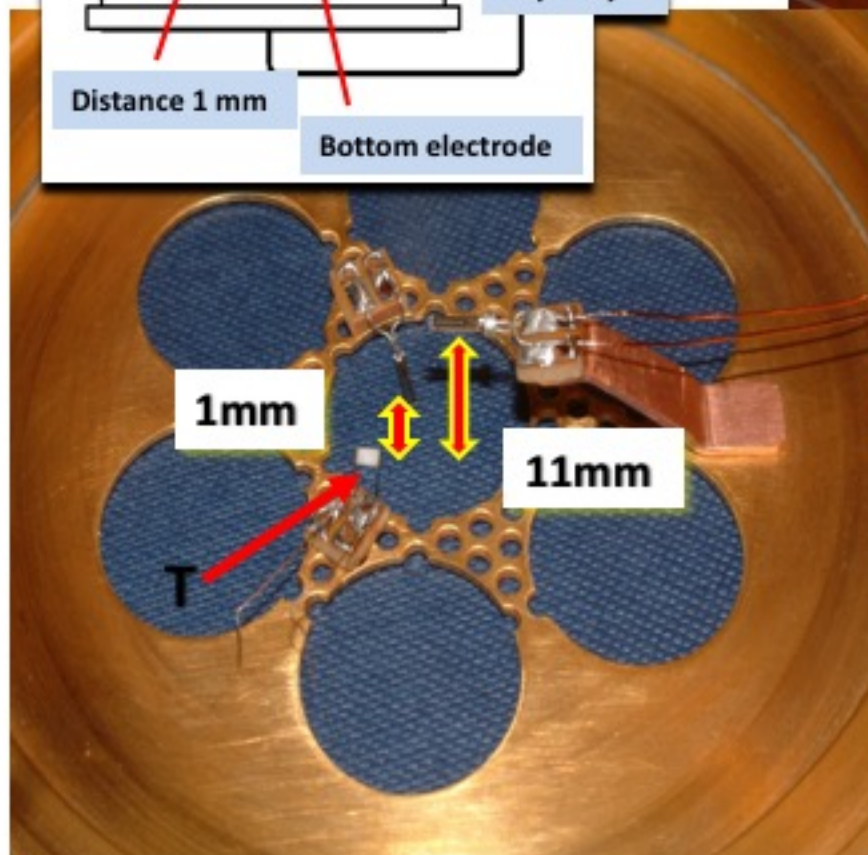
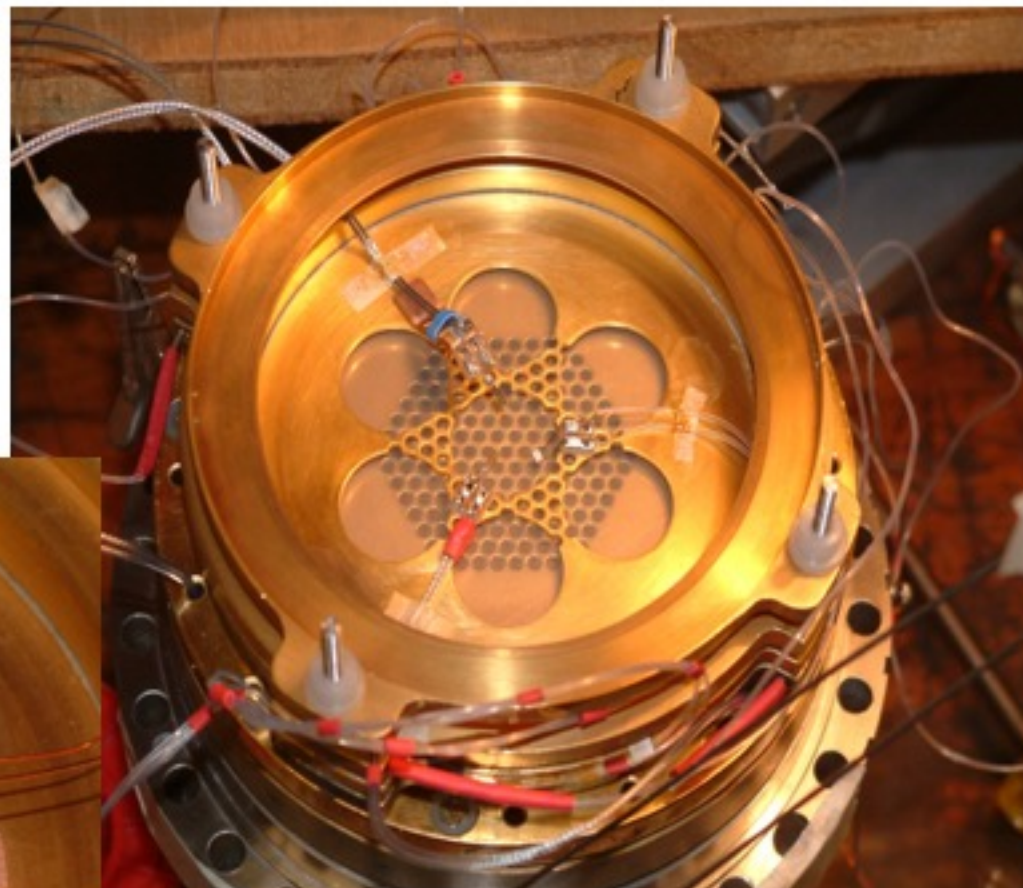
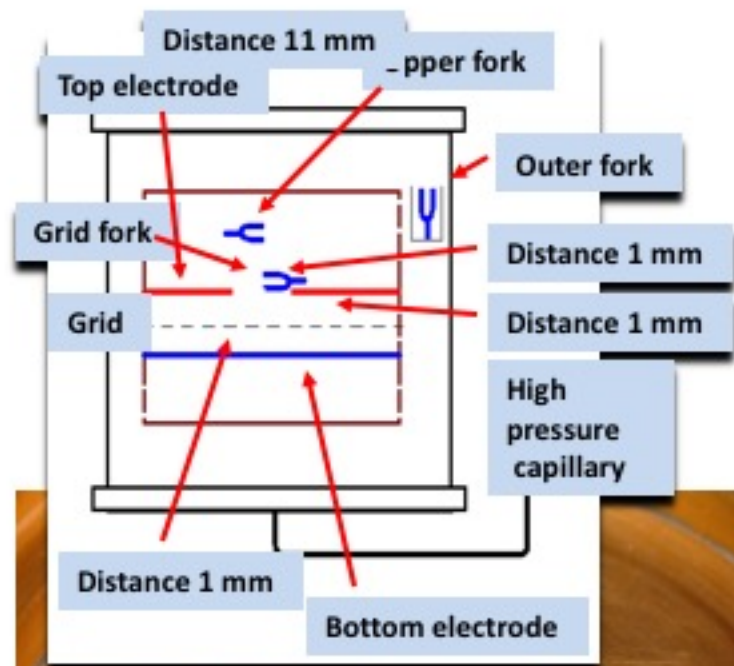
Vortex propagation in superfluid helium

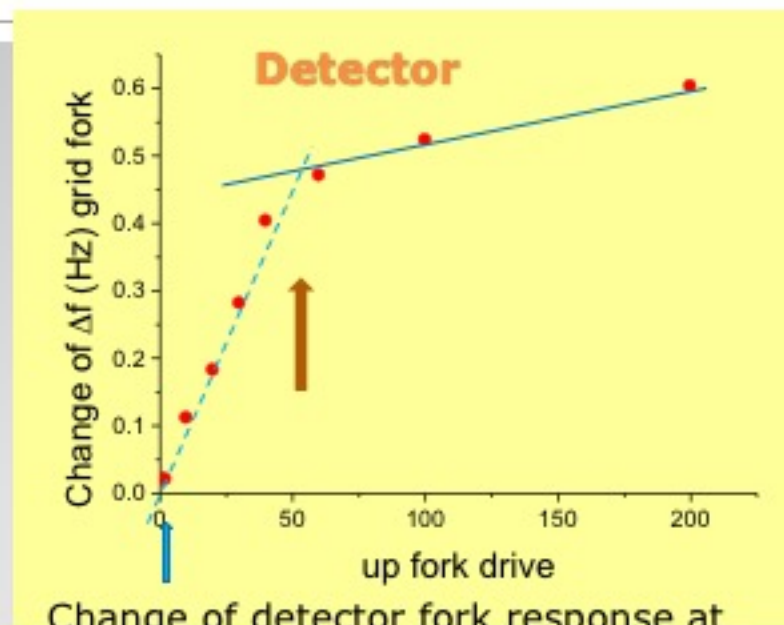
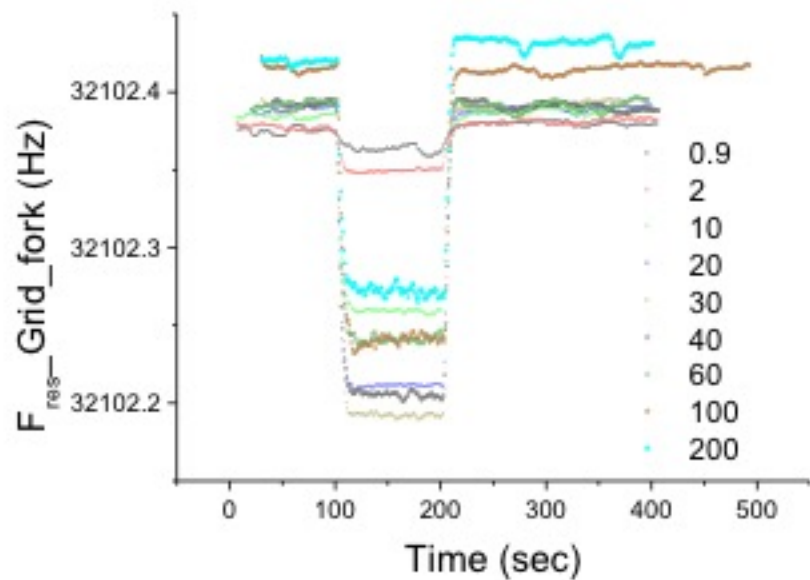
The have studied the vortex motion through superfluid helium. One of the standard 32 kHz fork was used as a generator of vorticity, another one was used as a detector of vortices. The scheme of experiment is shown on Fig.3.

The resonance frequencies of forks differ on 0.5 kHz.

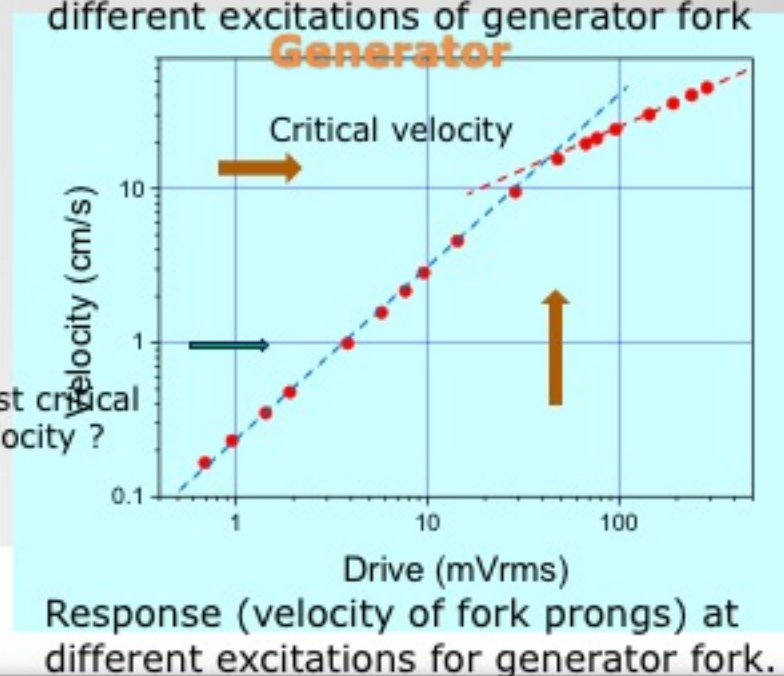
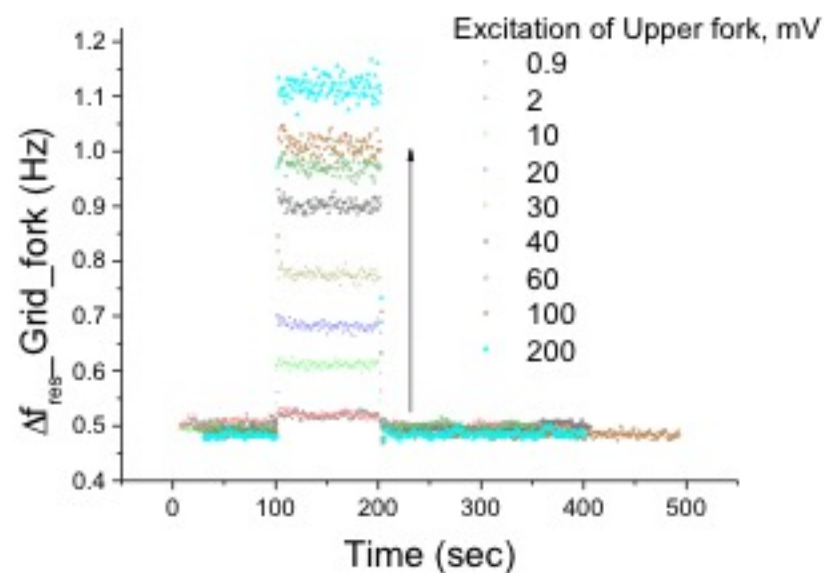
Fig 3. Scheme of experiments of vortex propagation. One of the forks was used as a generator with different excitation, another – as a detector with low oscillation. We mutual changed a function generator-detector for both forks.







Change of detector fork response at different excitations of generator fork



Response (velocity of fork prongs) at different excitations for generator fork.

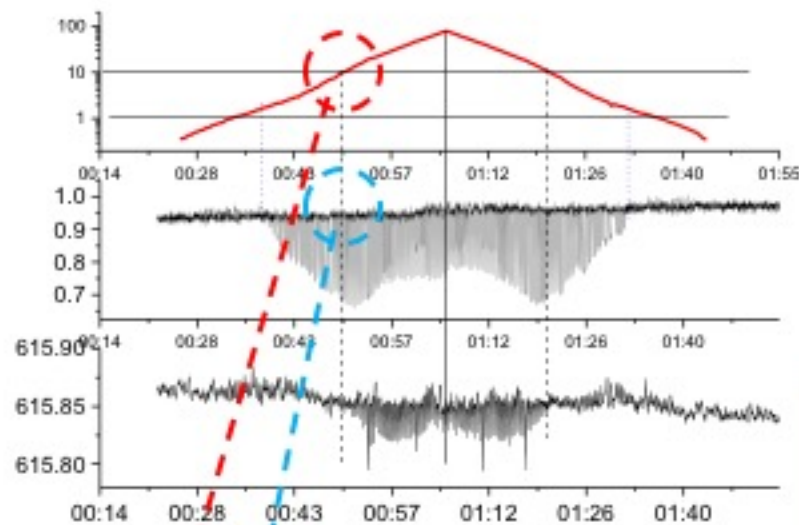
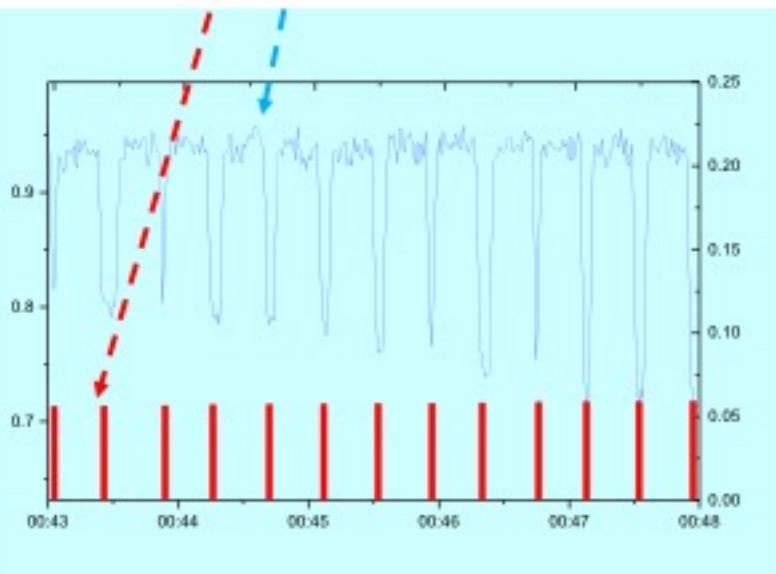


Fig 6: Time dependence of detector response at changing of generator excitation.

Excitation was done by telegraph code – signal with amplitude U_G - 2 s and waiting time with zero excitation 20 s.



Response (amp tracking at resonance point) of up fork with and without exciting the grid fork at fixed velocities shown in the box.

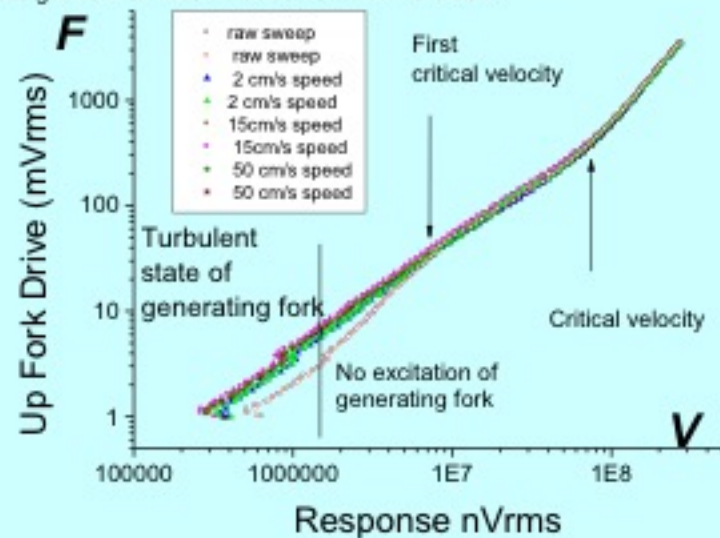
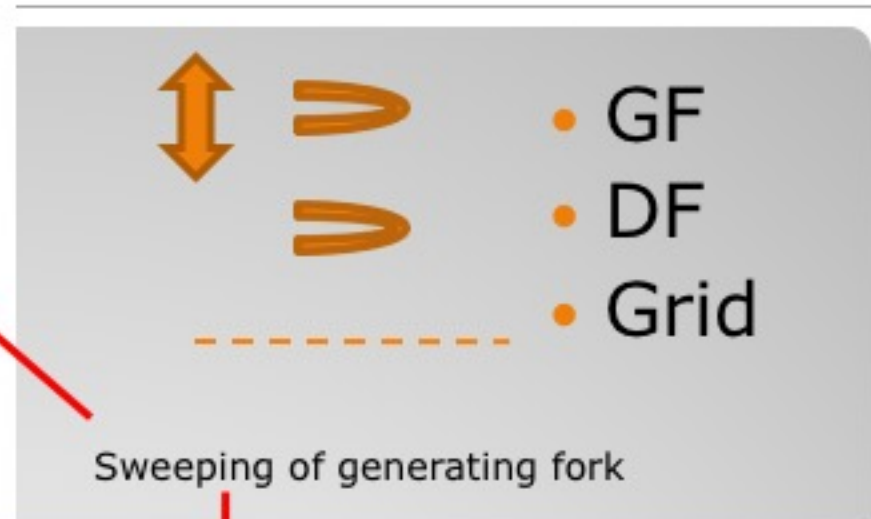
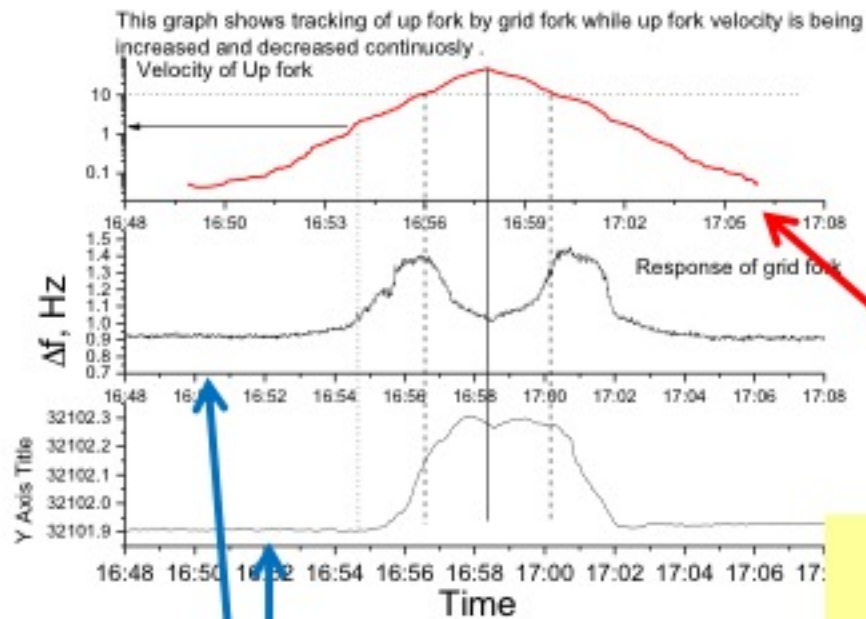
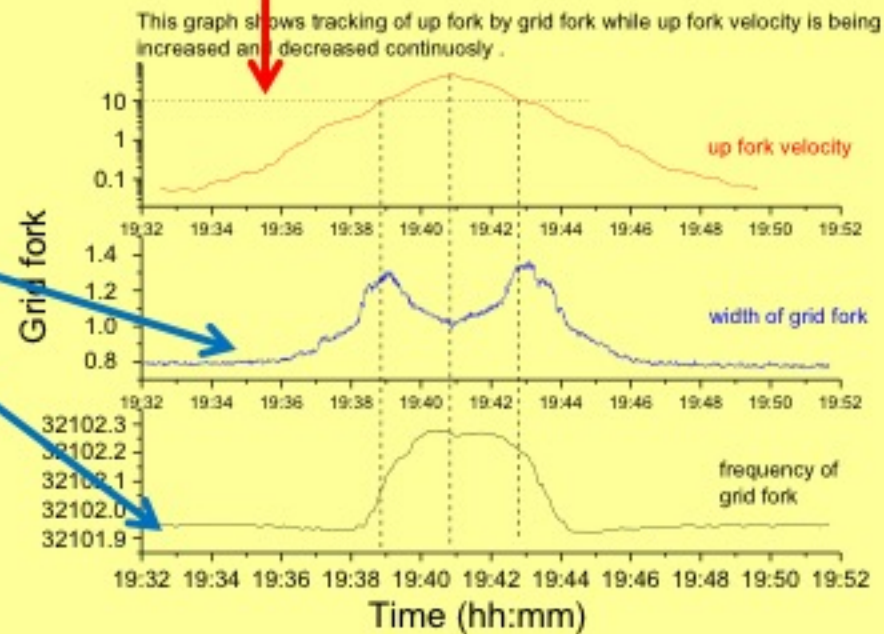


Fig 7: Influence of detecting fork response (prongs velocity) at different excitation (force of excitation and fluid resistance). The different colour corresponds to different fixed velocities of generating fork (and level of generated by this fork density of vortexes).

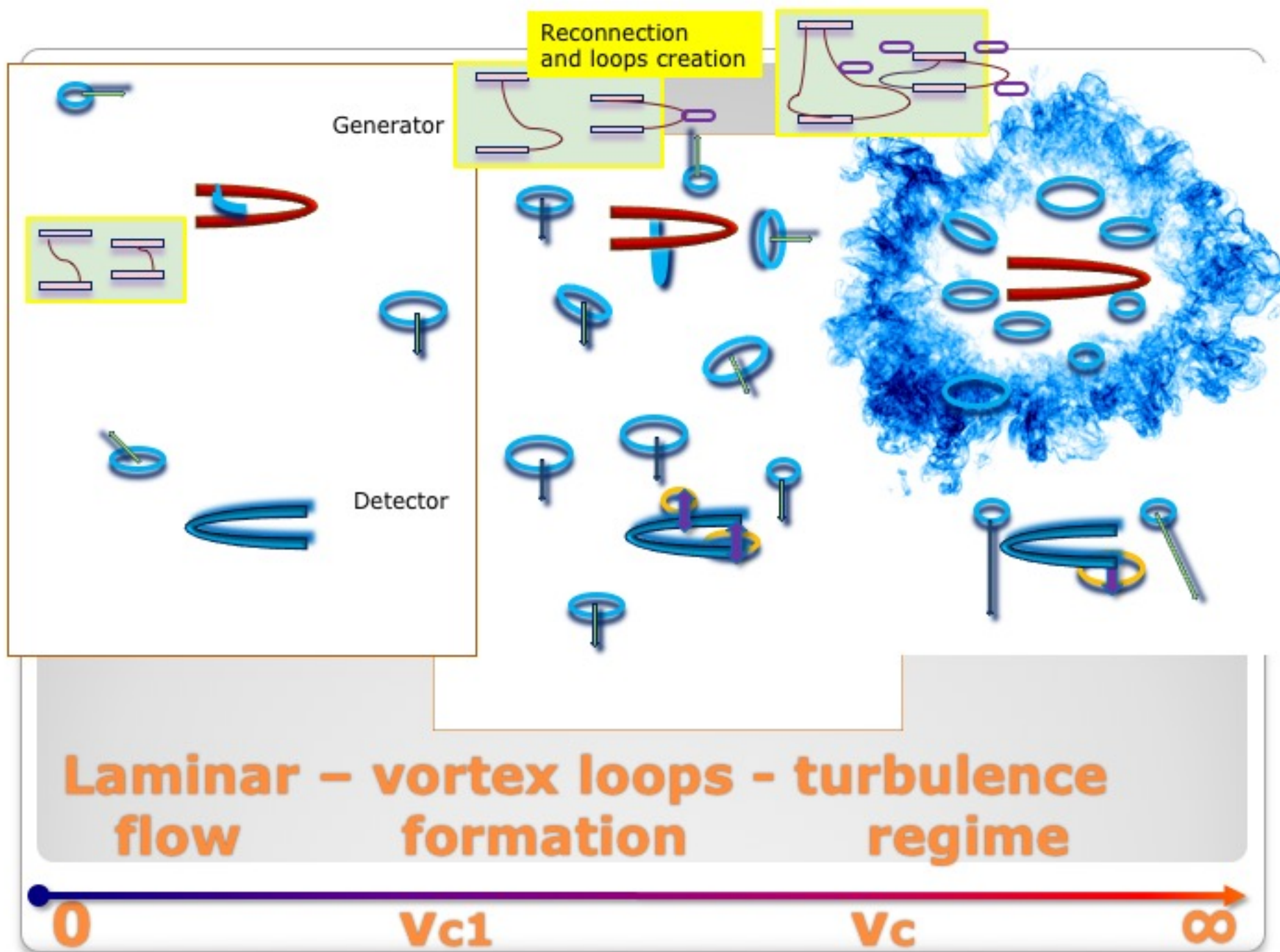


Response of detecting fork

Sweeping of generating fork



Tuning fork behaviour at low temperatures



Grid experiments

$$f_{R1} = 765 \text{ Hz}$$



[0,1]

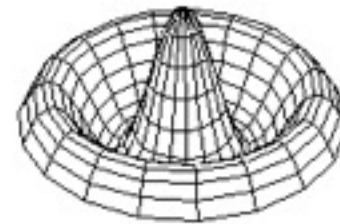
$$f_{T1} = \frac{2.405}{2\pi R} \sqrt{T/\sigma}$$

$$f_{R1} = 1756 \text{ Hz}$$



[0,2]

$$f_{T2} = 2.296 \cdot f_{T1}$$



[0,3]

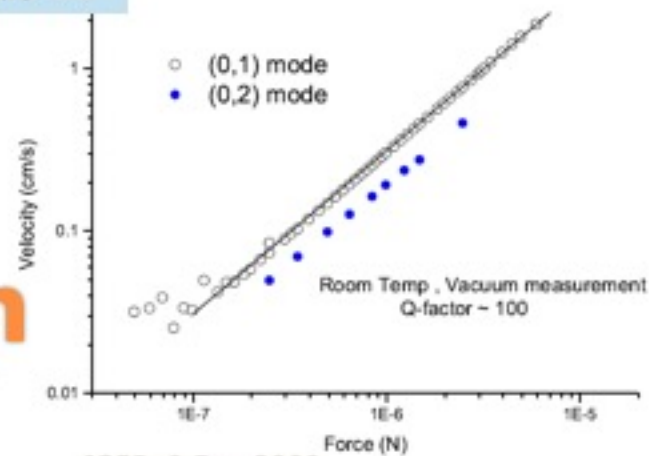
$$f_{T2} = 3.600 \cdot f_{T1}$$

T – surface tension, σ – area density, R – radius of membrane

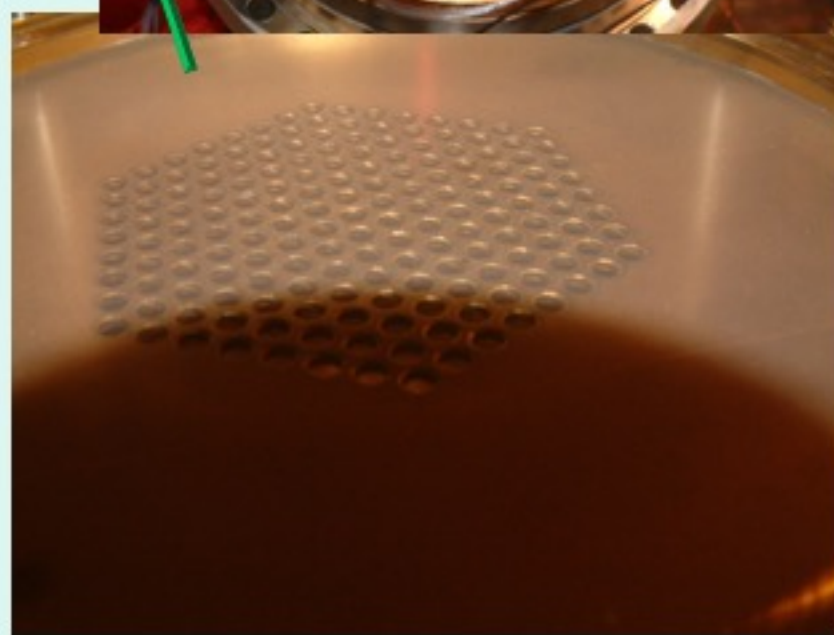
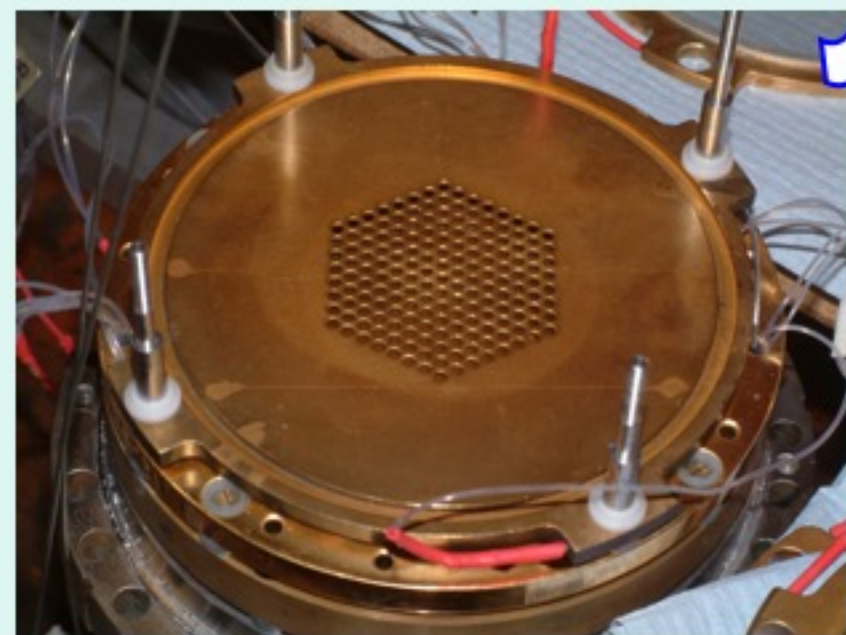
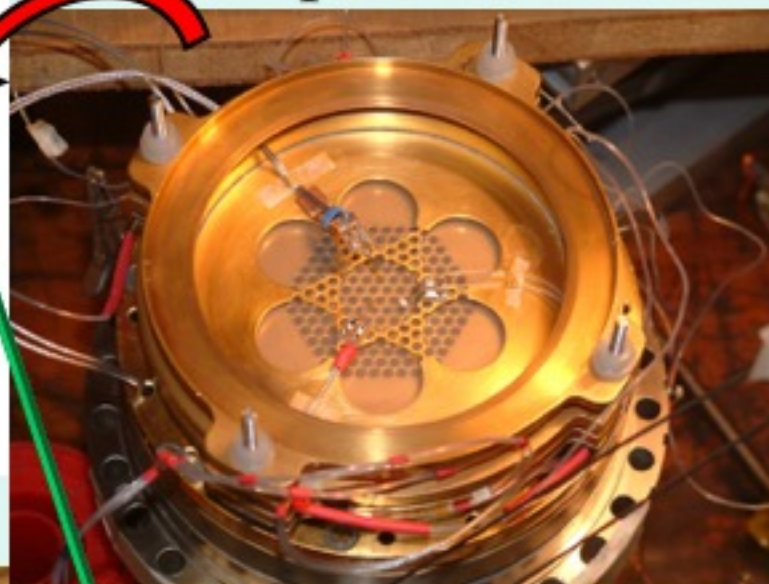
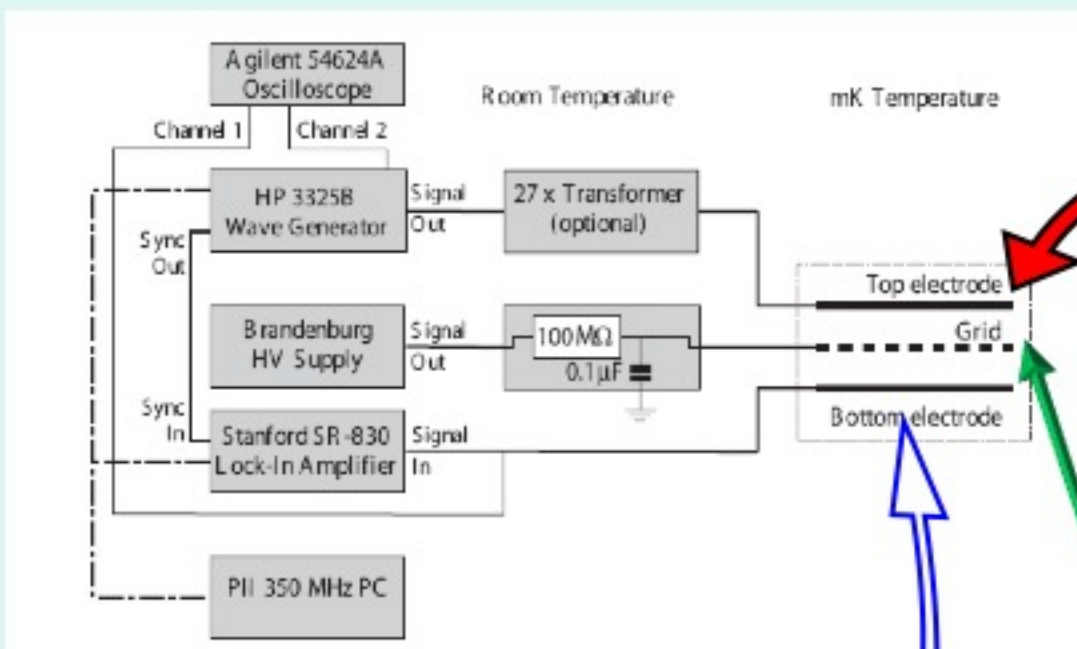
At room T $\Delta f = 7.2$ Hz, at base T (~ 10 mK) – $\Delta f = 0.004$ - 0.006 Hz

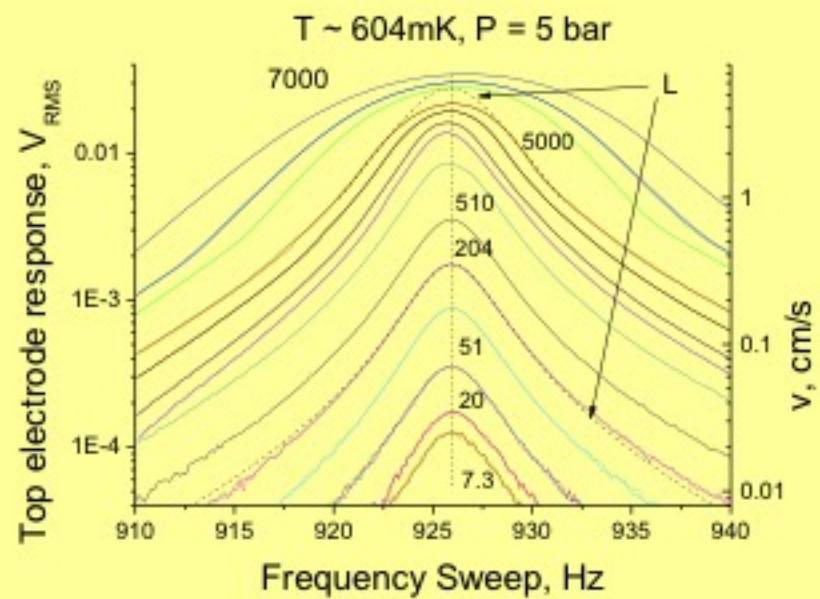
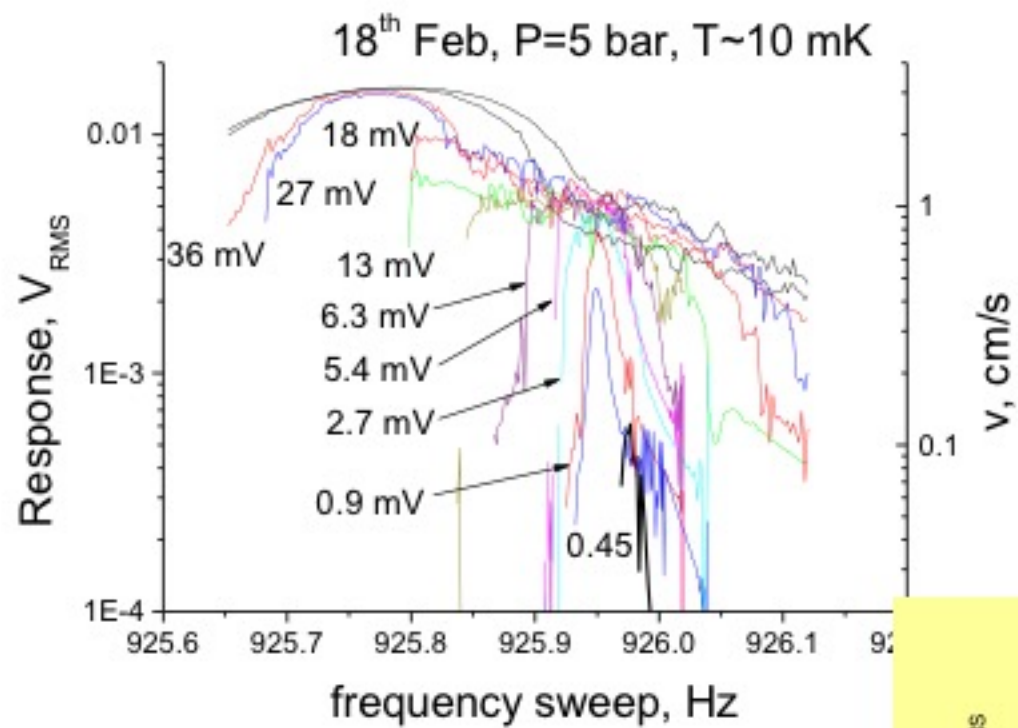
At room T $f \sim 765$ Hz, at base T (~ 10 mK) – $f \sim 925$ Hz

Membrane vibration

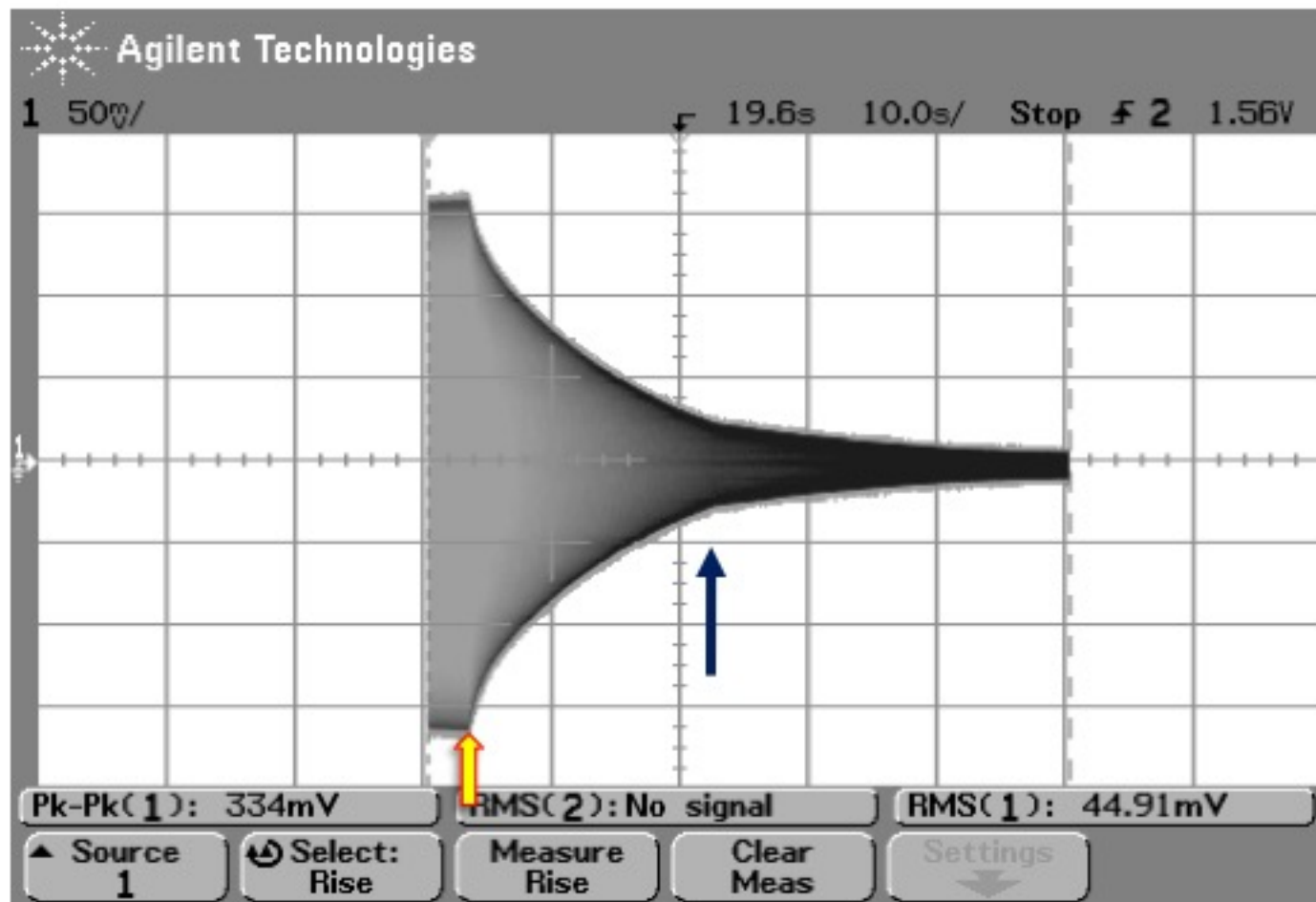


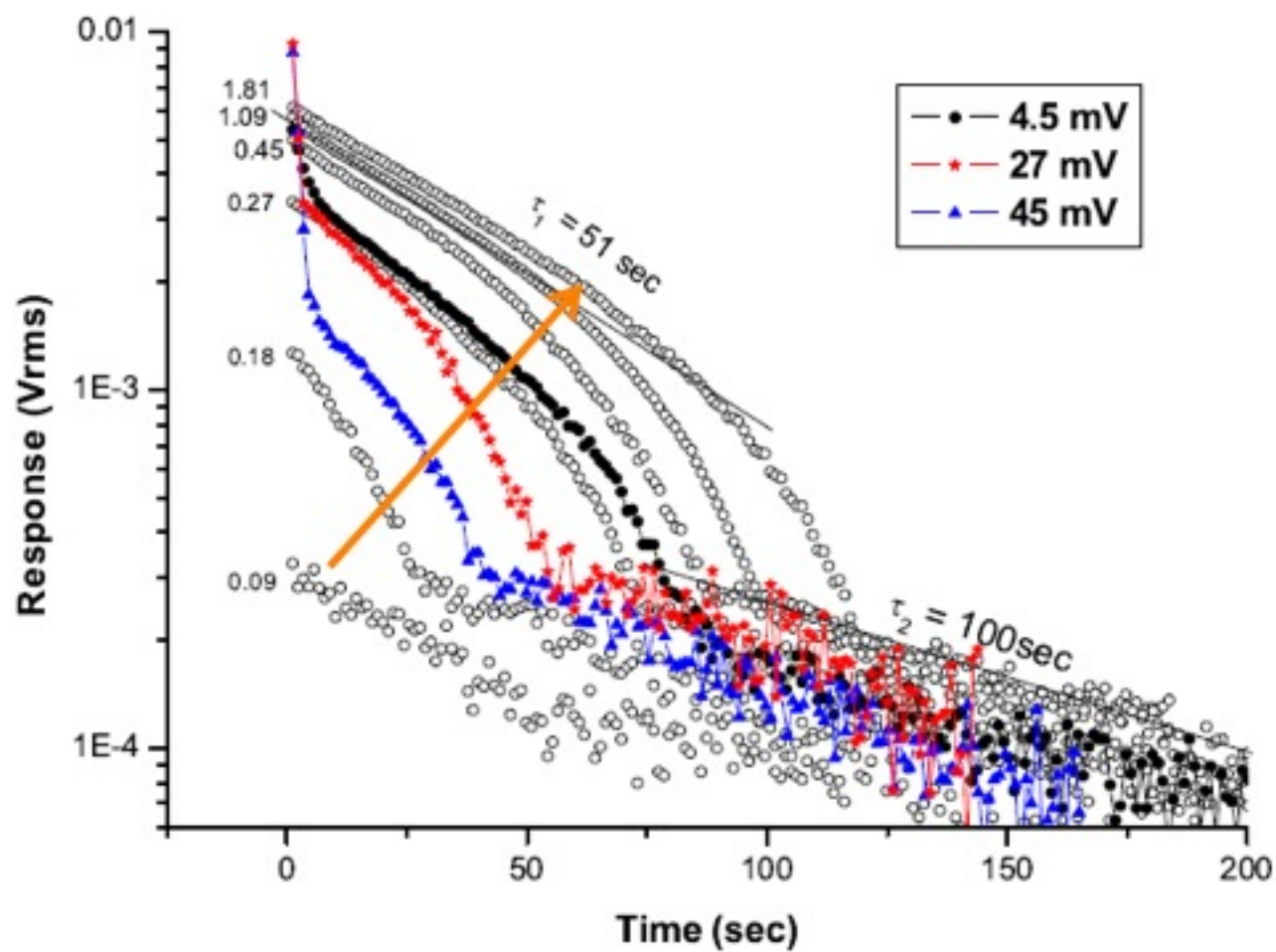
Scheme of experiments





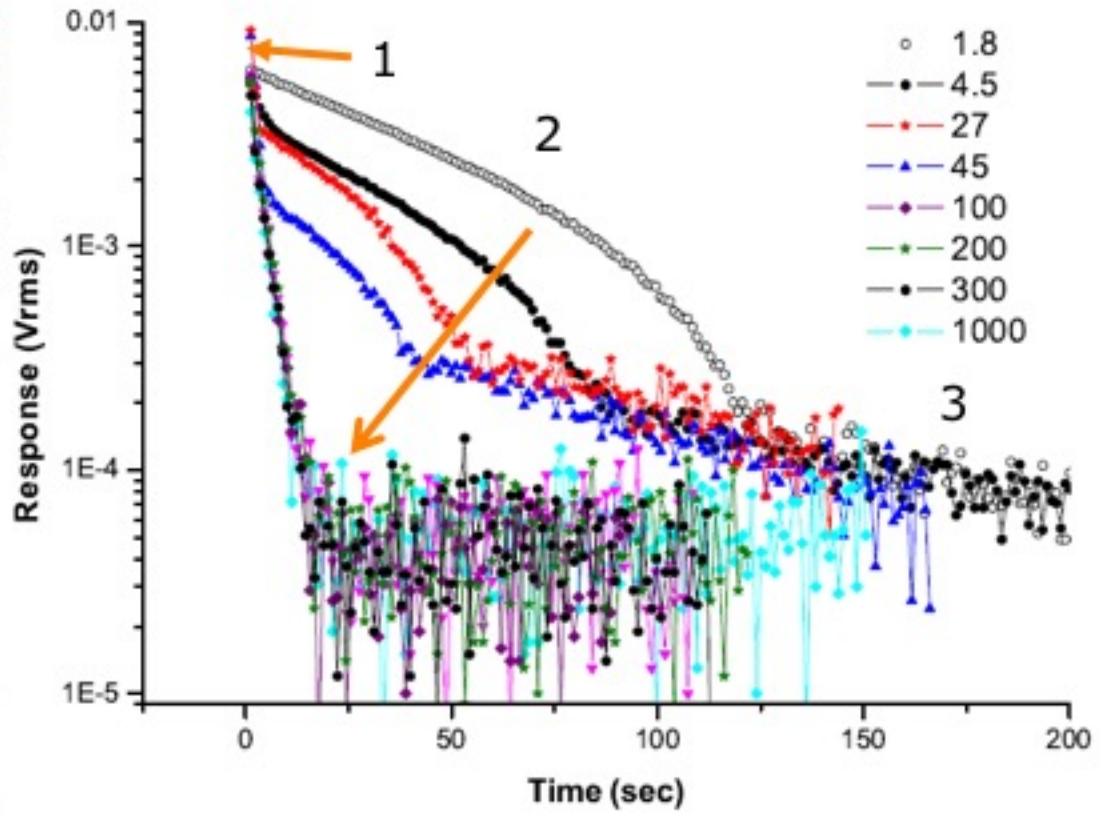
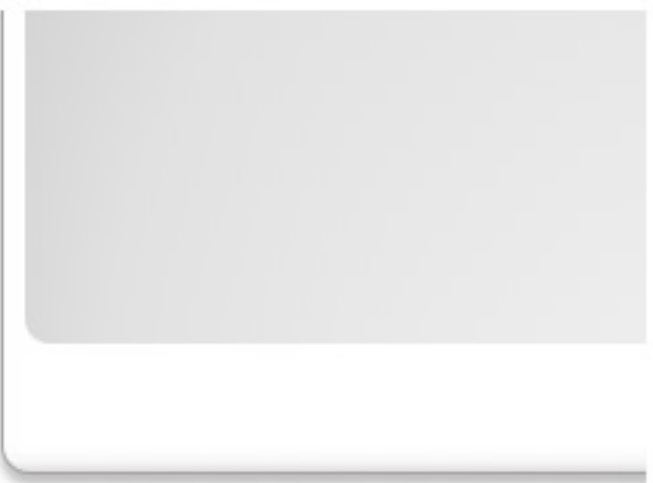
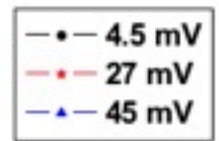
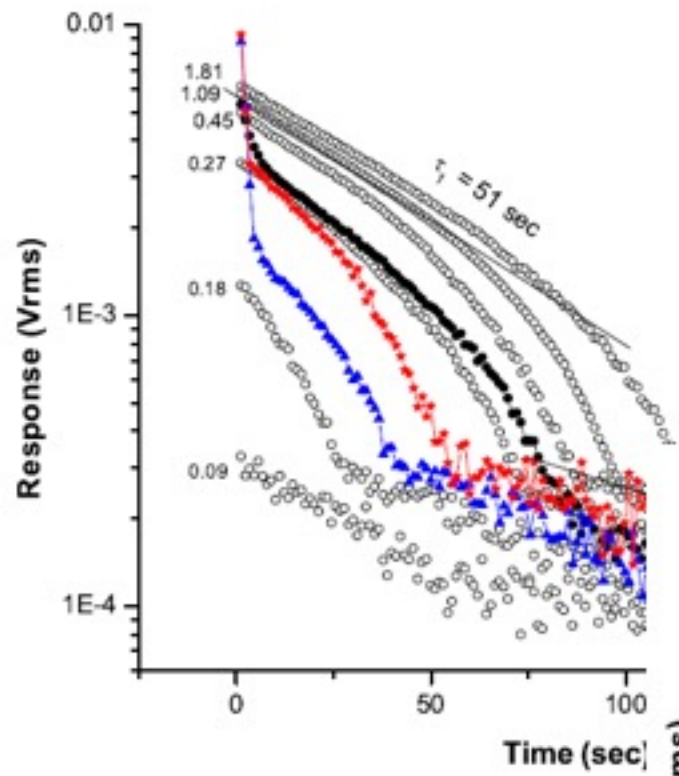
Free decay of grid oscillation





$T \sim 44 \text{ mK}$

Free decay of grid vibration



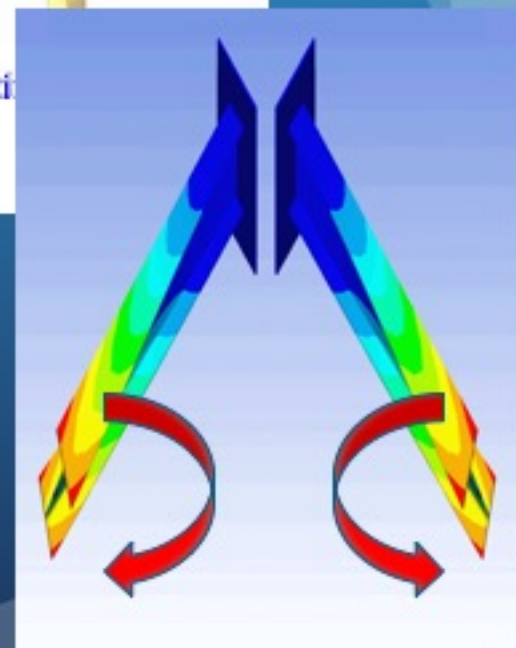
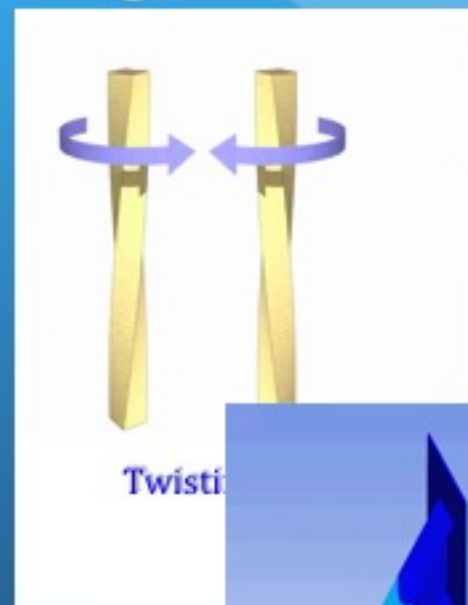


Torsion tuning forks

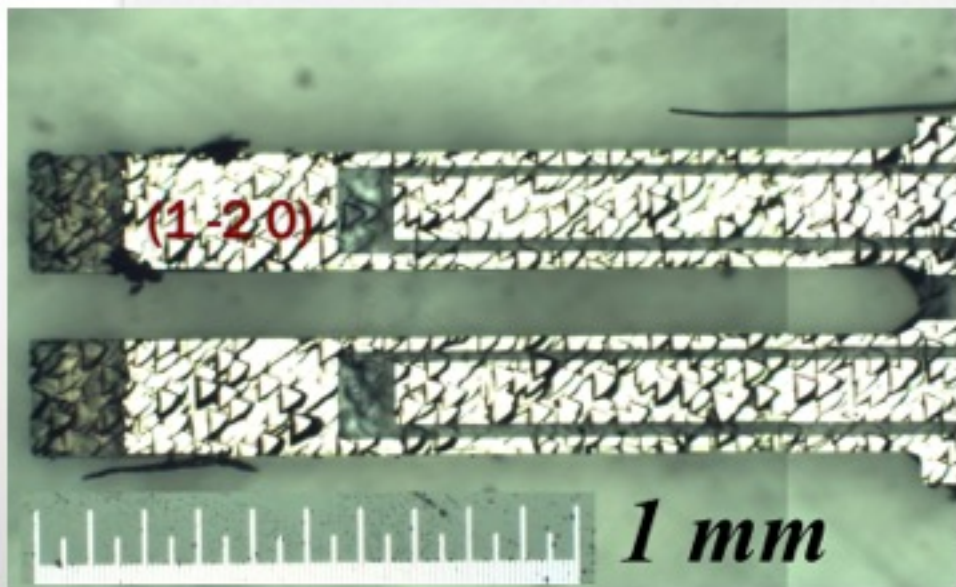
New tuning fork vibration modes

Flexible modes

Twisting modes



Size and vibration modes



$L=1.61$ mm - length

$T=0.223$ mm - thickness

$W=0.143$ mm - width

Mode of vibration: bending

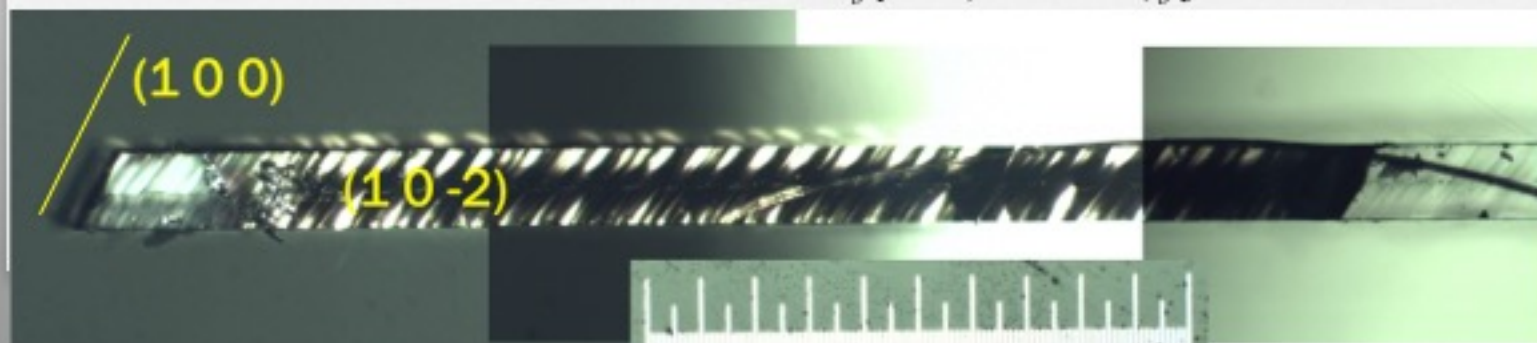
$$\omega_n = 2\pi f_n = \sqrt{\frac{EI}{\mu}} \frac{\beta_n^2}{L^2}$$

$f_1=77.2$ kHz, $f_2=482$ kHz

$f_3=1.35$ MHz

Measured vacuum T_{room}

$f_1=76,228$ kHz; $f_3=1.112$ MHz



Size and vibration modes

Mode of vibration: twisting

$$G = \left(\frac{2L\omega_n}{n}\right)^2 \rho_q K$$

$$K = \frac{\left(\frac{W}{T}\right) + \left(\frac{T}{W}\right)}{4\left(\frac{W}{T}\right) - 2.52\left(\frac{W}{T}\right)^2 + 0.21\left(\frac{W}{T}\right)^6}$$

$f_1=424$ kHz, $f_2= 848$ kHz

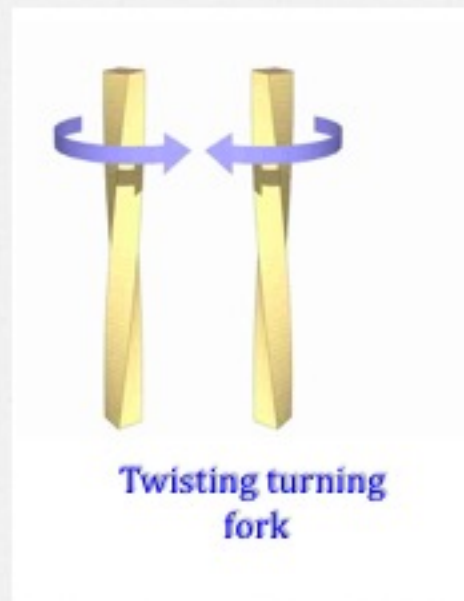
$f_3=1272$ kHz

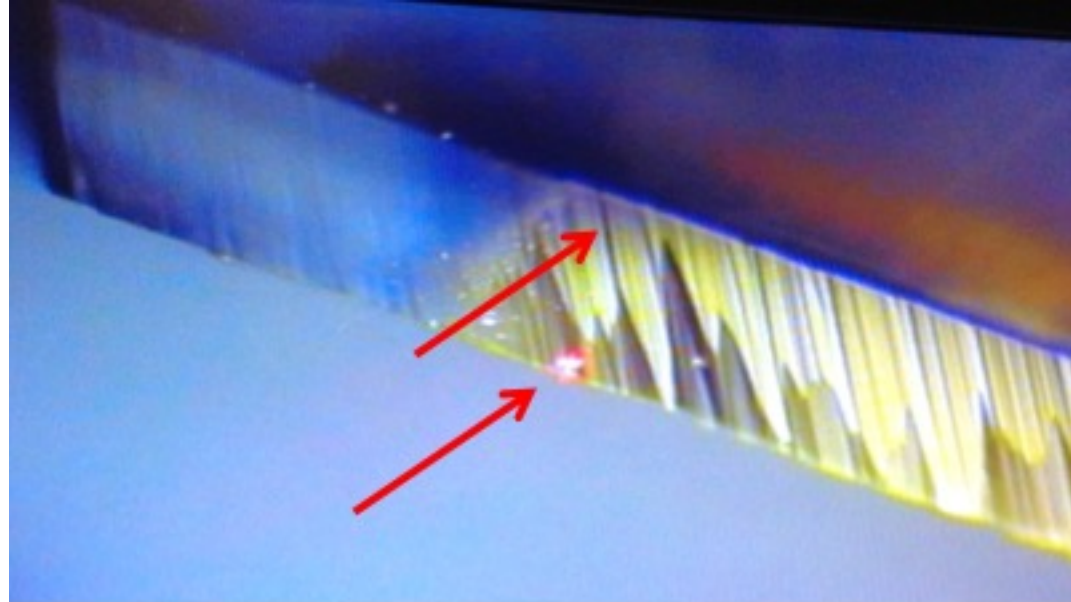
Measured vacuum T_{room}

$f_1=393\ 020$ Hz

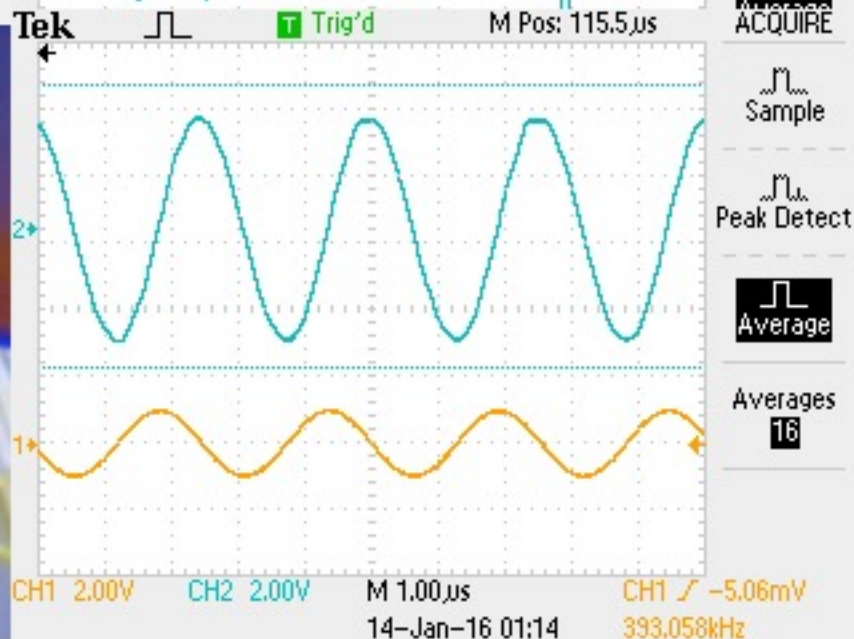
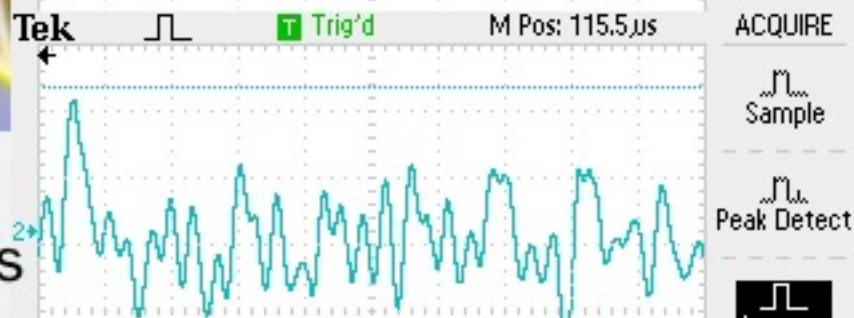
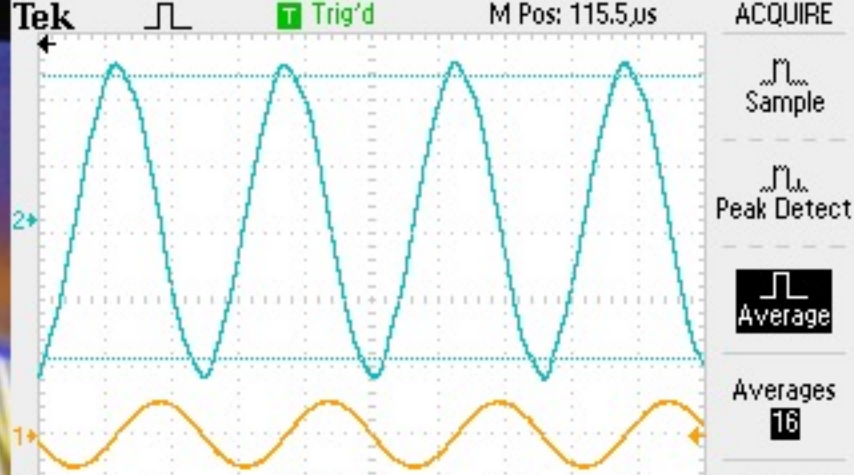
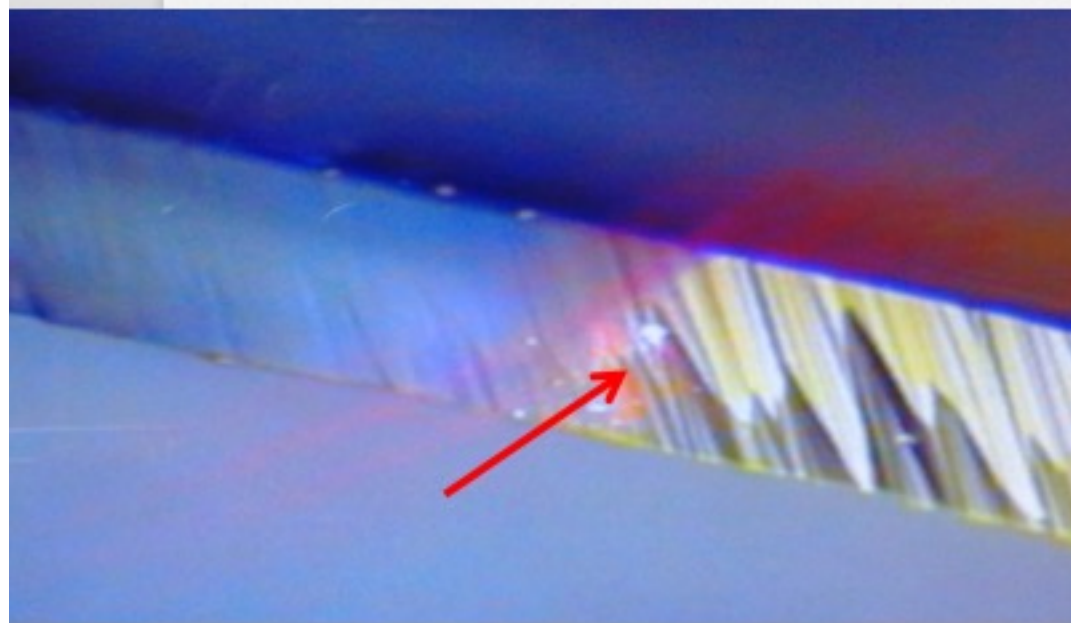
$$f_B \sim f_1 * (2n-1)$$

$$f_T \sim f_1 * n$$

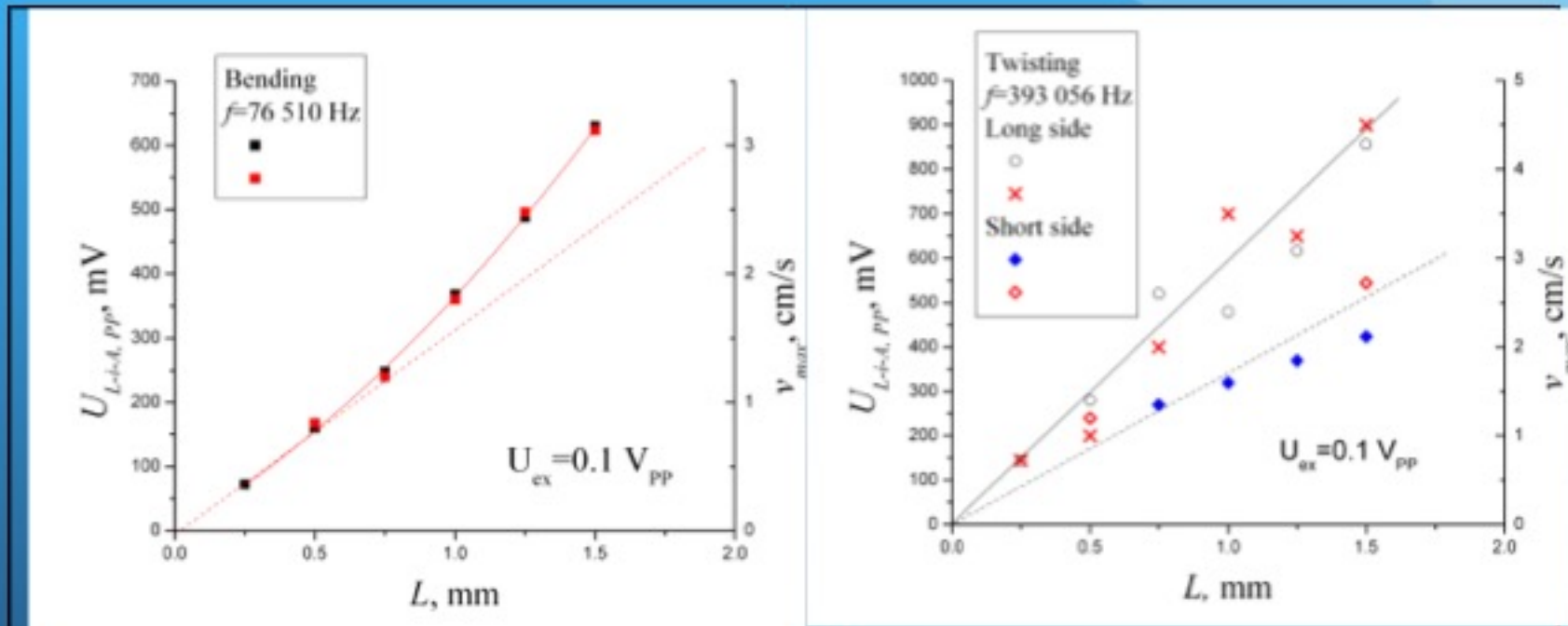




$f = 393\,093\text{ Hz}$, $\Delta f = 60\text{ Hz}$;
velocity of vibration $v_p = 10.4\text{ cm/s}$



Torsion fork motion at different modes



a)

b)

Fig.3 Scanning of the DA signal along the prong for different mode of vibration, a) bending mode and b) twisting mode from different sides of the prong. The right scale indicates a prong velocity (results 2017, air, T_{room}).

Different media of fork vibrations

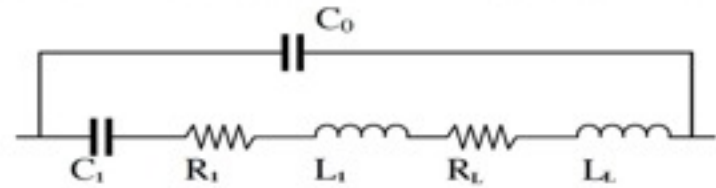
Calibration of I-V convertor.

393 kHz 0.571 V on R=33 kOm $U_{ex}=0$ dB
 76 kHz 0.629 V on R=33 kOm $U_{ex}=0$ dB

Frequency of resonance $U_{ex}=0$ dB (1V):

Gases	Vacuum Room	He Room Air	Room	N2 Low T
393kHz	393 020,1	393 014,75	392 973,5	392 653.5
76 kHz	75 663,24	75 688,8	75 643,75	75 927,6
Liquids	SF He T=1.4 K	He T=4.2 K	Liquid N ₂	N2 Low T
393kHz	387 892,5	388 372,5	371 997,5	392 653.5
76 kHz	74 308,9	74 458	67 567,5	75 927,6

Different media of fork vibrations

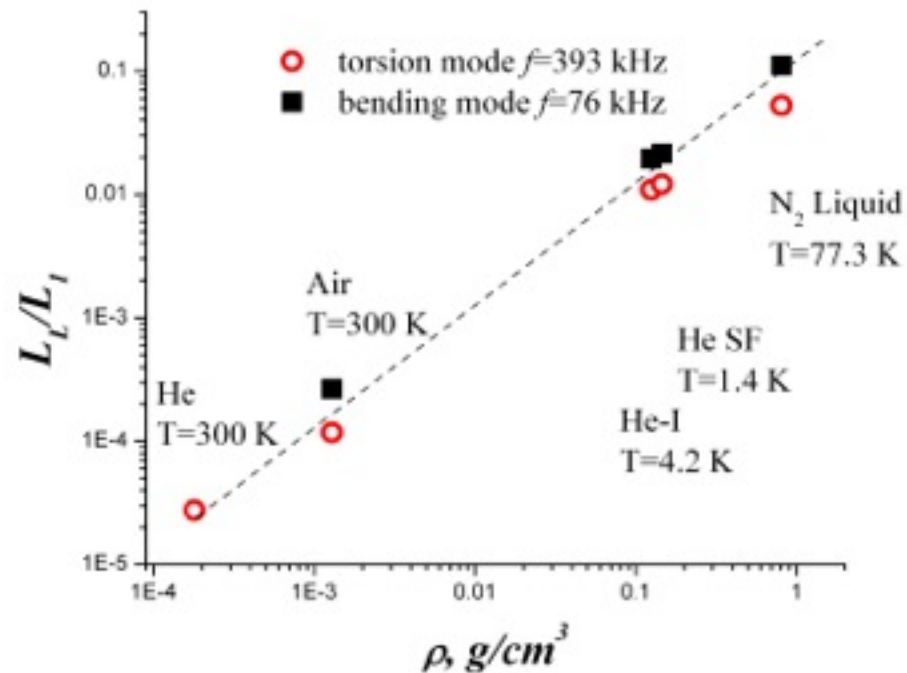


The resonant frequency of the circuit is given by

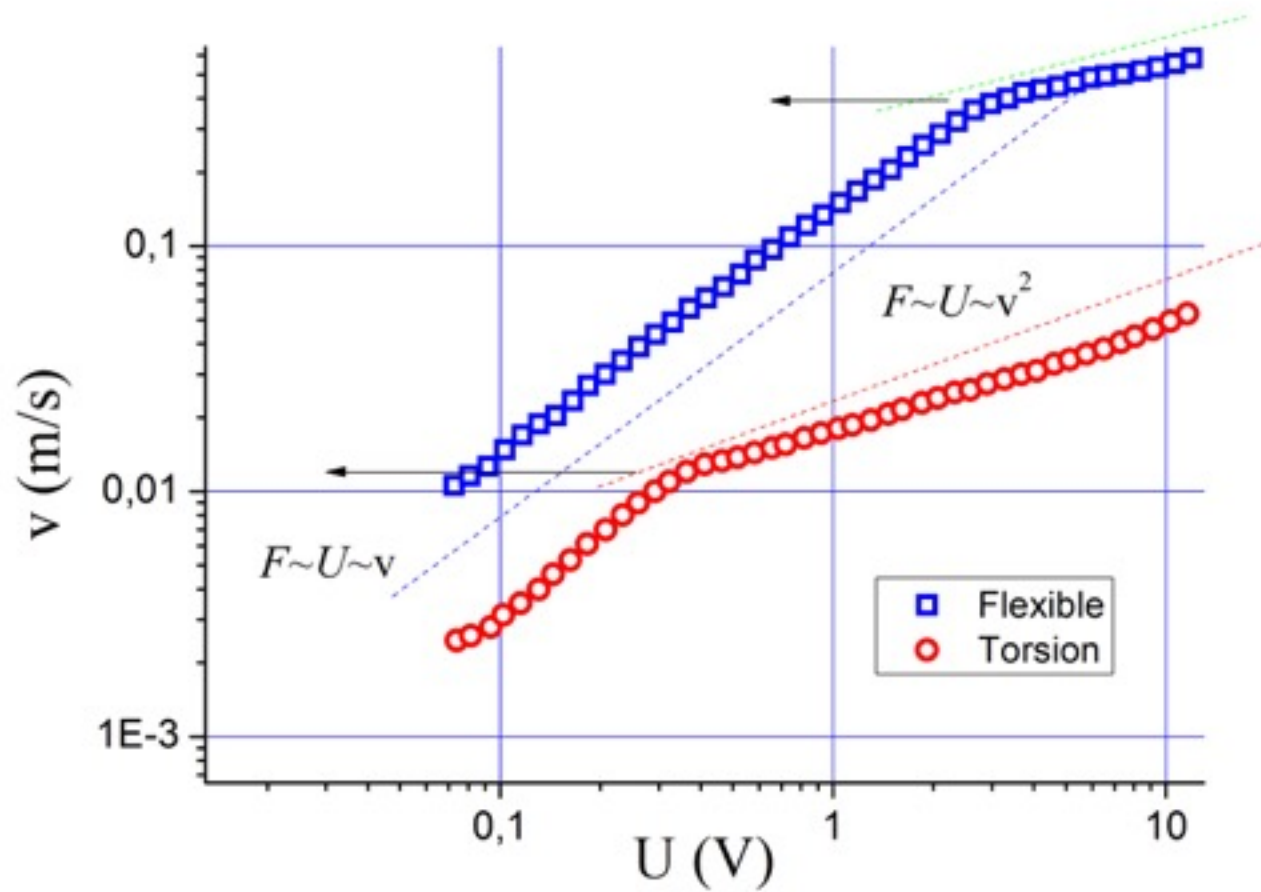
$$f = \frac{1}{2\pi\sqrt{(L_1 + L_L) * C_1}}$$

The resonant frequencies were measured in different media, so for these measurements changed only L_L

$$\Delta f \cong -\frac{1}{2} f_1 \frac{L_L}{L_1}$$



Fork vibration in superfluid helium



- Tuning forks are cheap, robust, reproducible, easy to install and very sensitive;
- Vibrating forks in superfluid helium use as a generator and a detector of turbulence;
- Influence of one fork on another at distance more bigger the size of fork – the vertexes move through cm's distance at $T \sim 0$ K;
- Virgin and disturb states exist in superfluid helium. The disturb state maintains hundreds hours at base temperatures ($T \sim 10$ mK);
- The energy loss of vibrating grid at free decay and tuning forks comes through three stages:
 - very quick energy loss – energy pumping into turbulent state (inject of vortex loop flux or classical turbulence?),
 - quick linear loss of energy (create of vortexes?),
 - slow free vibration (laminar motion, quality of grid);
- The torsion tuning fork has two modes: flexible and twisting. The transition into turbulent state has drastically different critical velocities for bending and twisting motion.

Conclusions