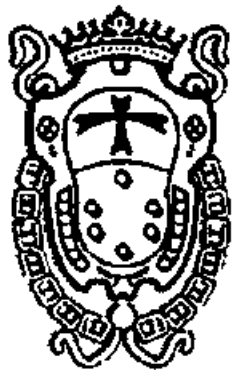


Decoherence in a Cooper pair Shuttle

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NATIONAL ENTERPRISE FOR NANOSCIENCE AND NANOTECHNOLOGY

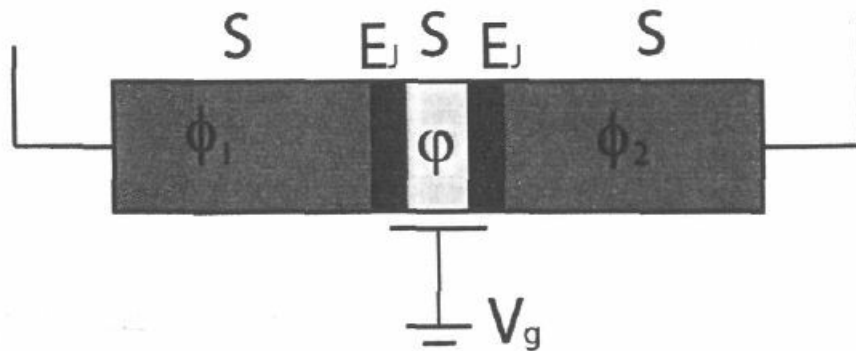
Scuola Normale Superiore - Pisa

Outline

- **Josephson vs Coulomb Blockade**
 - **Nanomechanical systems and shuttles**
 - **Cooper pair shuttle**
 - **Decoherence**
-

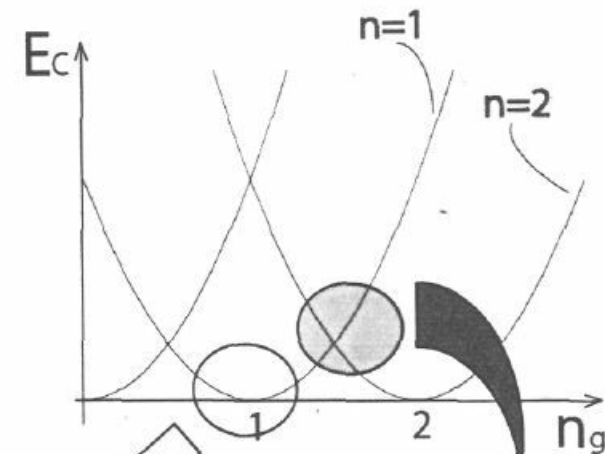
Josephson current in a S-SET

$$I = I_J \sin(\phi_1 - \phi_2)$$



The critical current
depends on the charging
energy of the island

$$H_c = E_c (\hat{n} - n_g)^2$$

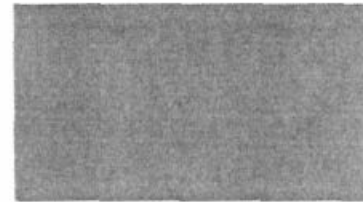
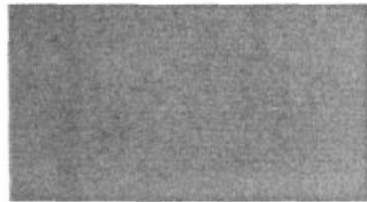


Coulomb
blockade

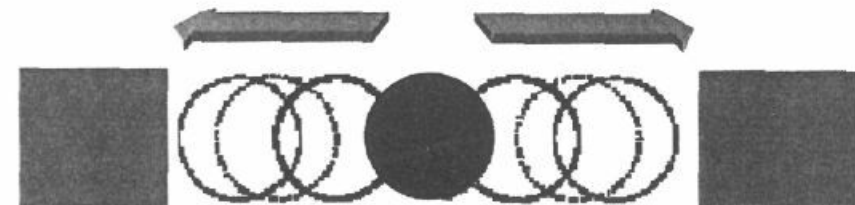
Degeneracy

Single Electron effects in the presence of mechanical vibrations

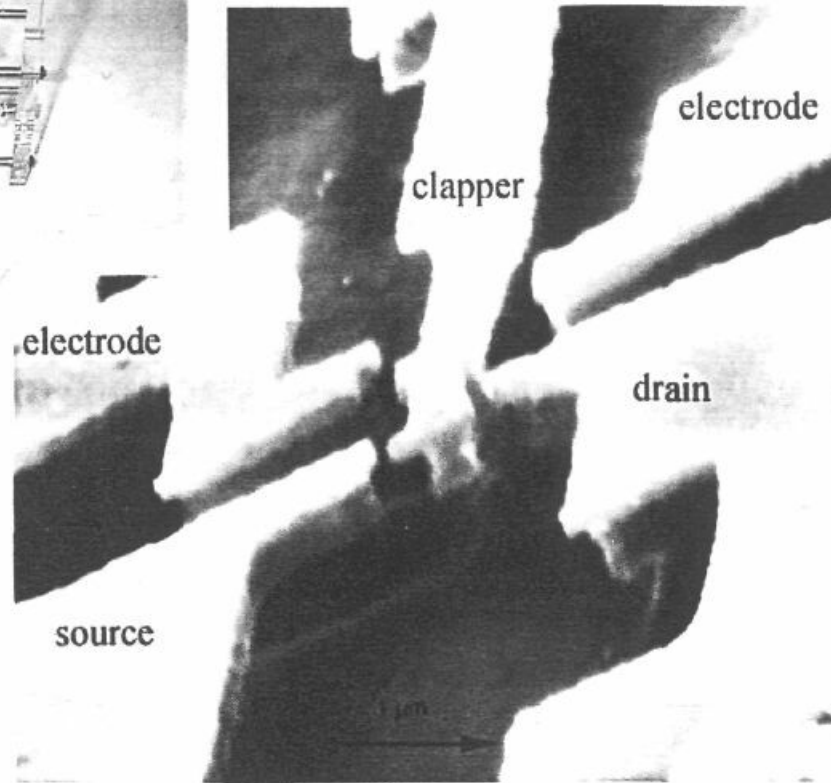
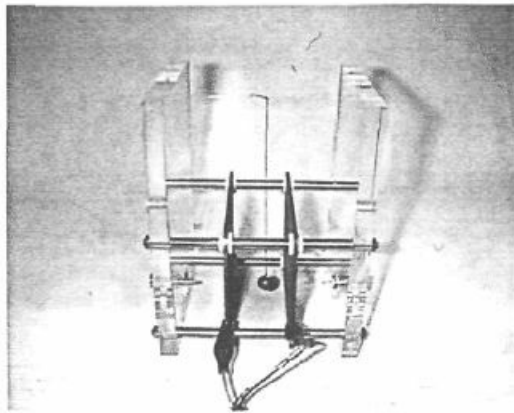
Gorelik et al 1998



SHUTTLE EFFECT



Experiments



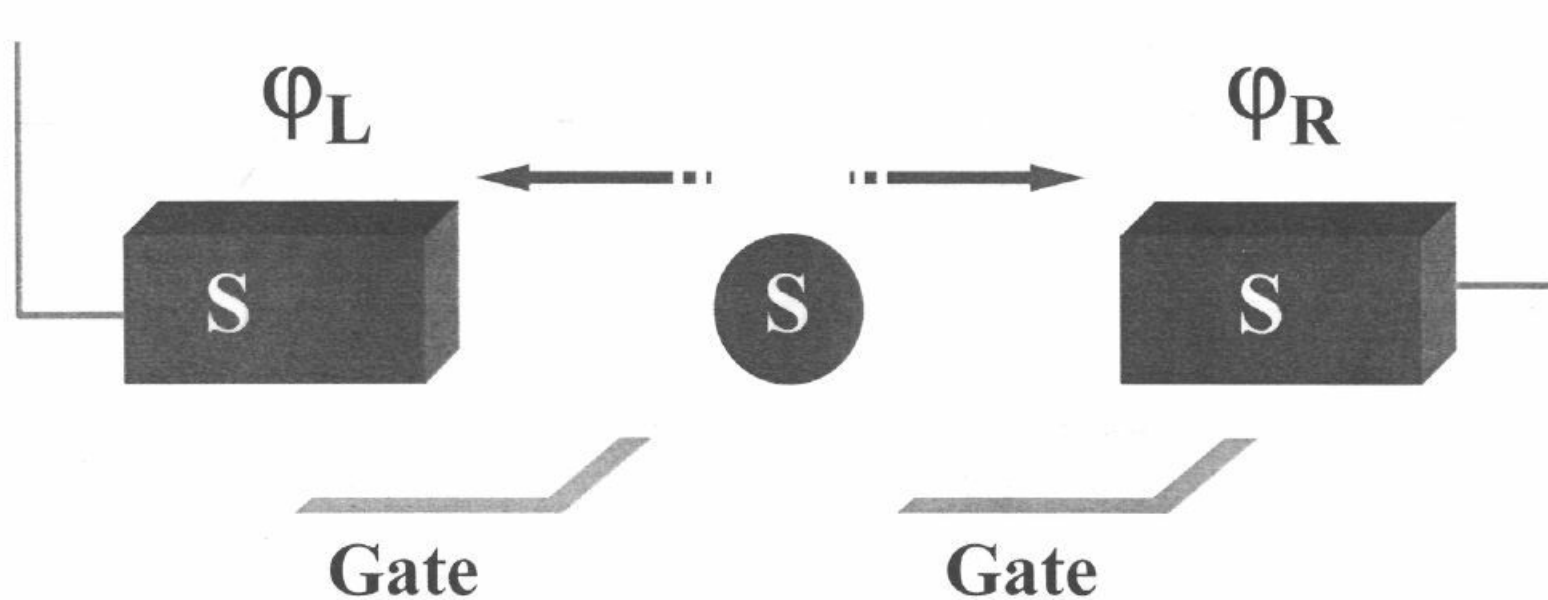
- Tuominen et al 1999
- Park et al 2000
- Erbe et al 2001

Superconducting Electrodes

- Shuttle effect with Cooper pairs?

Gorelik et al 2001

- Is it a coherent effect?
- What is the role of decoherence?



A single Cooper-pair box, by periodically moving between two superconducting leads, is able to keep phase coherence of the two distant electrodes

The Model

$$H = E_C (n - n_g(t))^2 - \sum_{\alpha=L,R} E_J^{\alpha}(t) \cos(\varphi - \varphi_{\alpha})$$

The external leads have well defined phase

- **The system is in regime of strong Coulomb blockade ($E_J \ll E_C$). The Hilbert space of the shuttle is two-dimensional, spanned by two charge states differing by one Cooper-pair.**



$$t_A \leq t \leq t_B$$

$$E_J^L \neq 0, E_J^R = 0, n_g = \frac{1}{2}$$

Free evolution



$$t_C \leq t \leq t_D$$

$$E_J^L = 0, E_J^R \neq 0, n_g = \frac{1}{2}$$



$$t_B \leq t \leq t_C$$

$$E_J^L = 0, E_J^R = 0, n_g = 0$$



“Josephson hybrid”

$$|\psi\rangle = \alpha |0\rangle + \beta e^{i\phi_{L/R}} |1\rangle$$



$$t_D \leq t \leq t_A + T$$

$$E_J^L = 0, E_J^R = 0, n_g = 0$$

Accumulated phases

- Phase difference

$$\bigcirc \varphi = \varphi_R - \varphi_L$$

- Dynamical phases

}

$$\bigcirc \vartheta = \frac{1}{2} \int_A^B E_J(t) dt$$

$$\bigcirc \chi = \frac{1}{2} \int_B^C E_C(t) dt$$

Decoherence

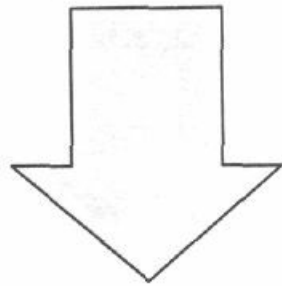
- Gate fluctuations
- Background charges
- Quasi particle tunneling
- ...

Important in establishing a stationary state

Strongly affects the Josephson current

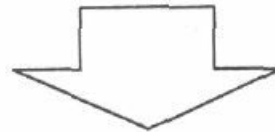
Coupling to the environment

$$H_{\text{coupling}} = \sum_i \lambda_i (b_i + b_i^+)$$



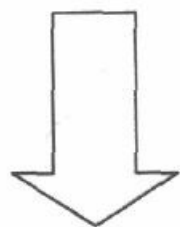
Treated in Born-Markov approximation

Master equation for the density matrix ρ in the space generated by $|0\rangle, |1\rangle$ states



Steady state : fixed point for the map $\rho(nT) \rightarrow \rho(nT + T)$

The effect of the environment are taken into account by



Contact region

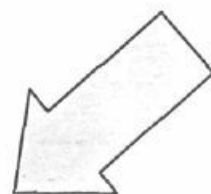


Free evolution region



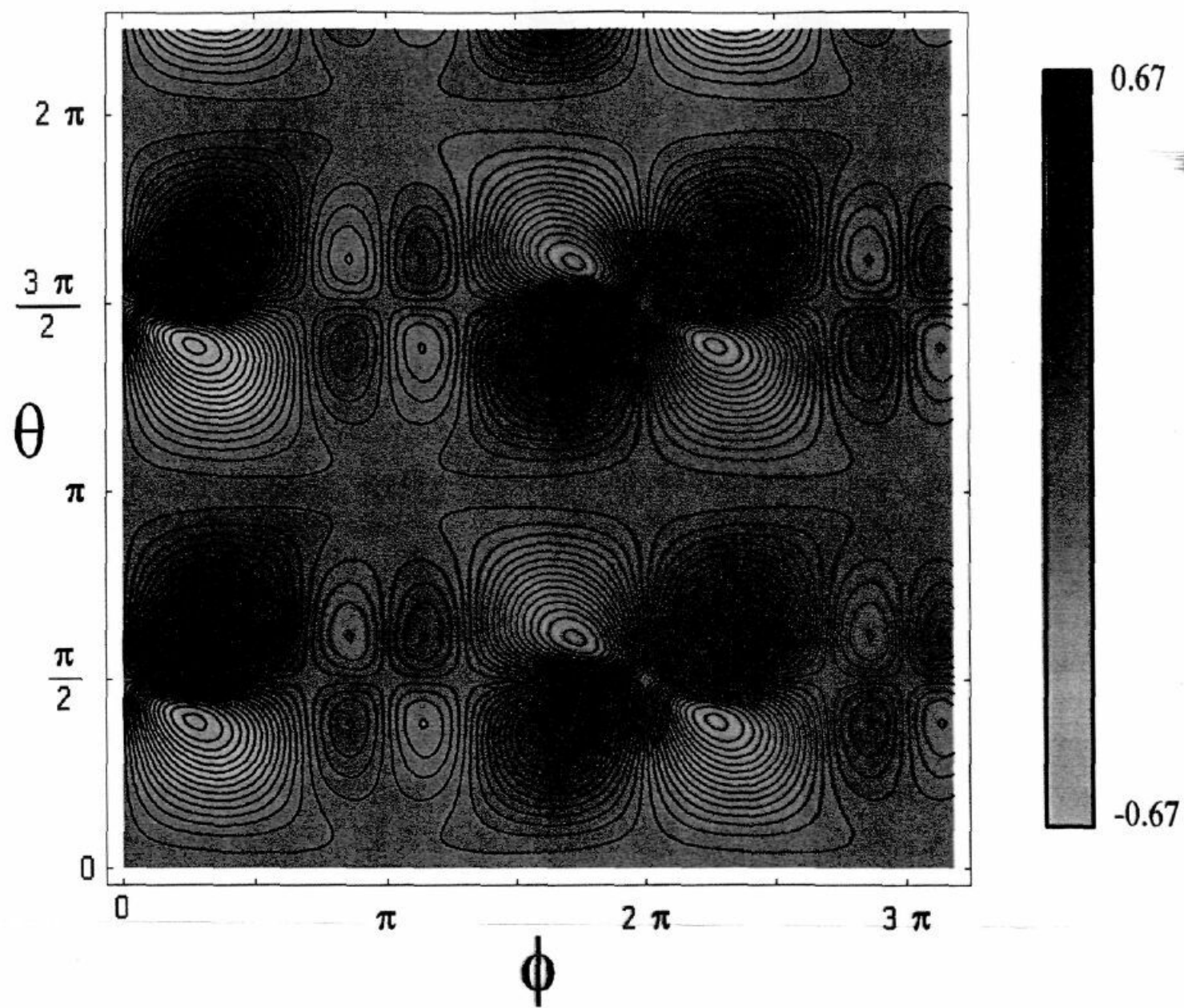
Modification of the critical current:

$$I = I(\phi, \mathcal{G}, \chi, \gamma_J t_J, \gamma_c t_c)$$



Average current over a period:

$$I = \frac{1}{T_{\text{period}}} \int \langle \hat{I}(t) \rangle dt$$



Weak damping

$$\gamma_J t_J \ll \gamma_C t_C \ll 1$$

$$I \approx \frac{2e}{T} \frac{\gamma_J t_J}{\gamma_C t_C} \frac{(\cos \varphi + \cos 2\chi) \tan \vartheta}{1 + \cos \varphi \cos 2\chi} \sin \varphi$$

Current is suppressed if dephasing in the contact region tends to zero

Strong damping

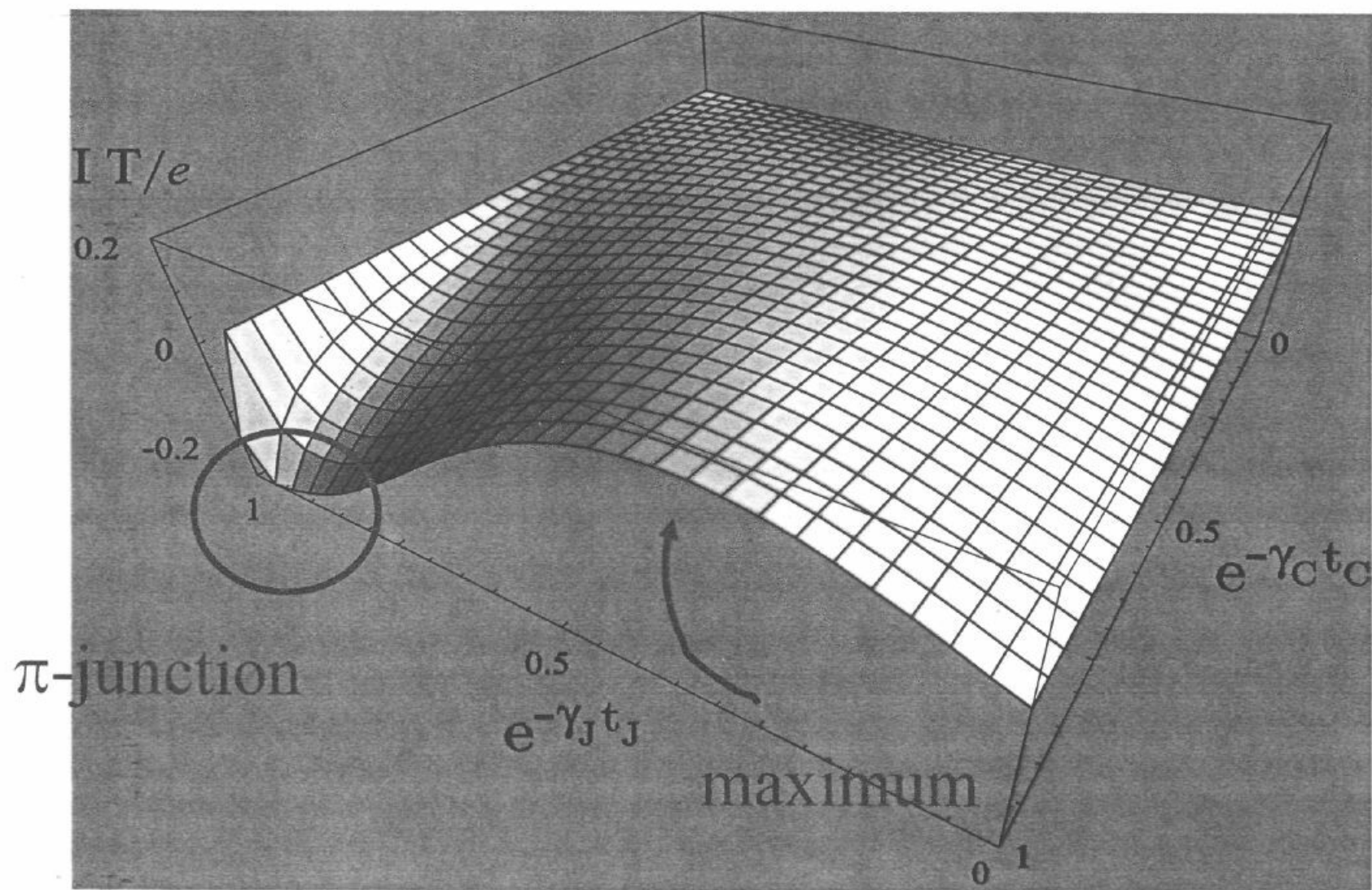
$$\gamma_J t_J \gg 1, \gamma_C t_C \gg 1$$

$$I = \frac{2e}{T} \tanh(2\beta E_J) e^{-(\gamma_J t_J + \gamma_J t_J)} \cos(2\chi) \sin(2\vartheta) \sin(\phi)$$

If the dephasing is strong enough
the coherence is lost before the
period is completed: the current
is exponentially suppressed

Change in sign

Dependence on the dephasing rates



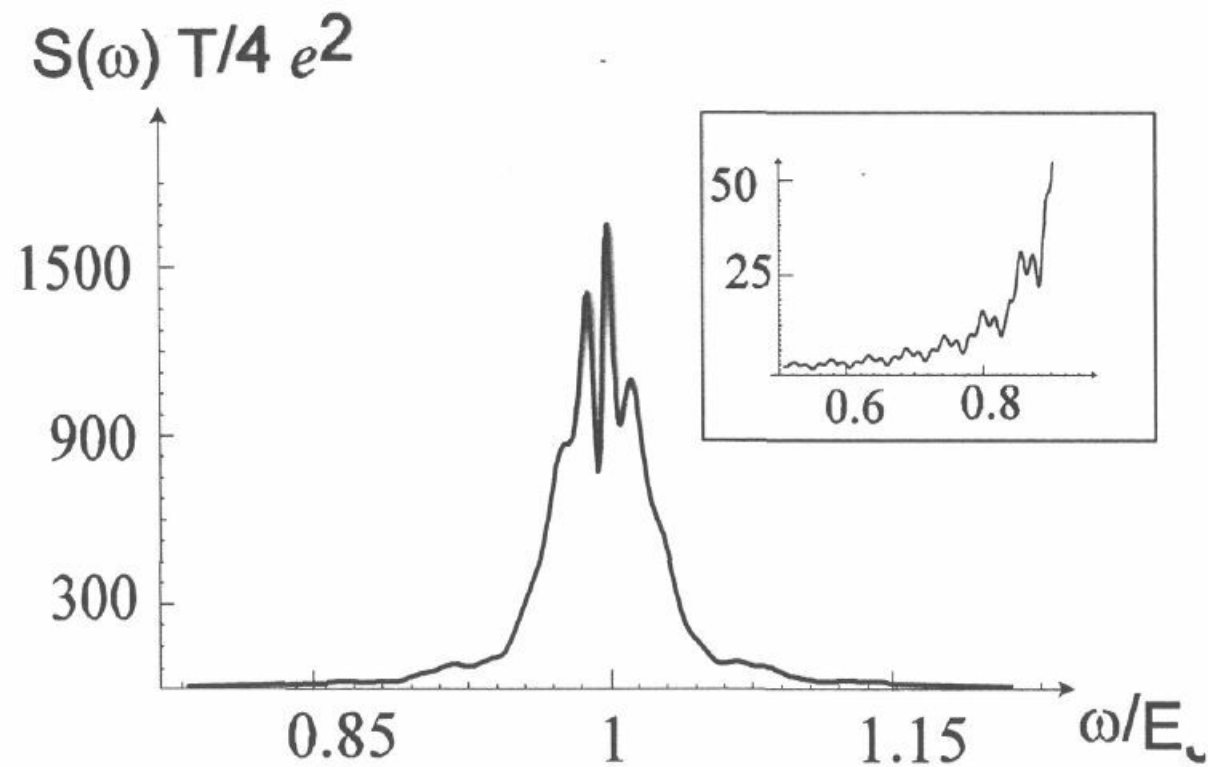
Zero-frequency noise

Strong damping $S(0) = \frac{4 e^2}{T} \left\{ \frac{1}{2} - e^{-\gamma_J t_J} \cos \theta + e^{-2\gamma_J t_J} f(\vartheta, \chi, \phi, \dots) \right\}$

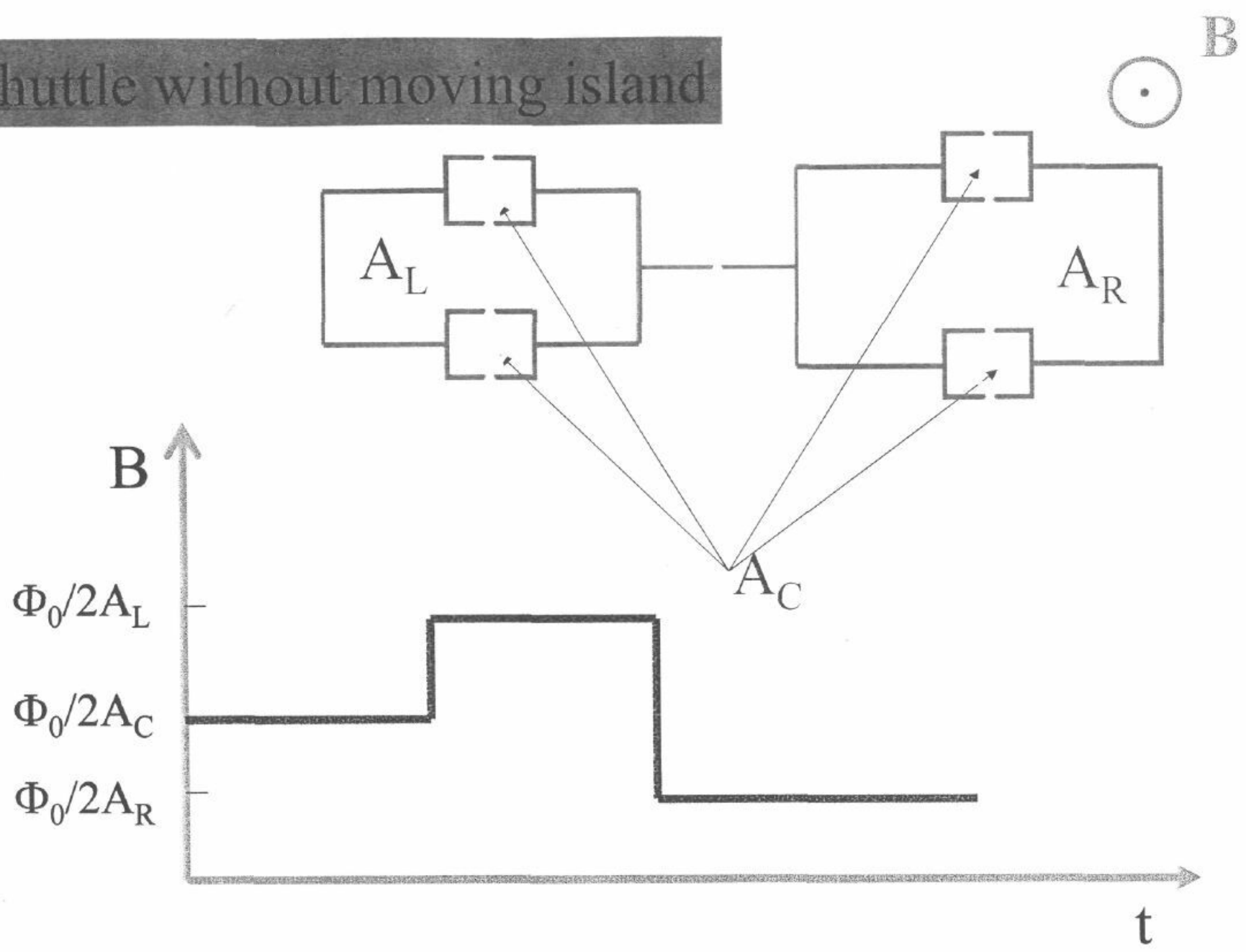
Weak damping $S(0) = \frac{4 e^2}{T} \frac{1}{\gamma_c t_c} \frac{\tan^2 \vartheta \sin^2 \phi}{2(1 + \cos \phi \cos(2\chi))}$

Current noise

$$S(\omega) = \frac{1}{T} \int_0^T dt \int_{-\infty}^{+\infty} d\tau \left\{ \frac{1}{2} \langle \hat{I}(t+\tau) \hat{I}(t) \rangle - \langle \hat{I}(t+\tau) \rangle \langle \hat{I}(t) \rangle \right\} e^{-i\omega\tau}$$



Shuttle without moving island



Conclusions

- The Cooper pair shuttle is able to maintain and create phase coherence between two “distant” superconductors.
- Fluctuations of gate voltage can either enhance or suppress the supercurrent.
- The Cooper pair shuttle can ~~have~~ ^{behave as} a π -junction behavior.
- The current noise displays a peak at Josephson energy.